VERTICAL FLUX OF TRACER IN THE GULF OF MEXICO: THE IMPACT OF SMALL SCALE DYNAMICS AND STORMS

- 3 Jouanno¹, J., X. Capet², J. Sheinbaum³, F. Durand¹, R. Dussurget⁴, and E. P. Chassignet⁵
- 4 1 LEGOS, Université de Toulouse, IRD, CNRS, CNES, UPS, Toulouse, France.
- 5 2 CNRS-IRD-Sorbonne Universités, UPMC, MNHN, LOCEAN Laboratory, Paris, France
- 6 3 Departamento de Oceanografía Física, CICESE, Ensenada, Baja California, Mexico
- 7 4 Mercator Ocean, Ramonville-Saint-Agne
- 8 5 Center for Ocean-Atmospheric Prediction Studies, Florida State University, Tallahassee, FL, USA
- 9 Corresponding author: Julien Jouanno (julien.jouanno@ird.fr)

10 Key Points:

- 11 High resolution simulations of the Gulf of Mexico have been developed
- Eddy activity on scales shorter than the mesoscale reduces the exchanges between the surface and the thermocline
- Restratifying action of the submesoscale turbulence limits the deepening of the mixed-layer

16 Abstract

17 In a series of regional simulations (with horizontal resolution at ¹/₄°, 1/12° and 1/36°), configured

18 to investigate the role played by the meso- and submesoscale turbulence on the dispersion of

- 19 passive tracers in the upper layers of the Gulf of Mexico, we show that eddy activity on scales
- 20 shorter than the mesoscale reduces the exchanges between the surface and the thermocline. This
- 21 is in contrast to previous studies which suggest that enhanced submesoscale activity in the winter
- 22 mixed layer may actually lead to increased exchanges with the permanent thermocline. This
- 23 reduction is explained by the restratifying action of the submesoscale turbulence that effectively
- 24 limits the deepening of the mixed-layer in response to atmospheric winter synoptic events.

25 1. Introduction

26 Submesoscale processes, including eddies and filaments at scale of O (10) km and below, have

- 27 received increasing attention in recent years, stimulated by their potential key role in the energy
- dissipation of the balanced flows (Capet et al. 2008, Capet et al. 2016), their impact on the large
- 29 scale circulation (Lévy et al. 2010) or their consequences on the biogeochemistry (Mahadevan and
- 30 Archer 2000, Lévy et al. 2001, Lévy et al. 2012). Submesoscale oceanic currents have multiple

origins, the most predominant being mixed-layer instability, strain induced frontogenesis or
 topographic wake (see the review by McWilliams 2016).

33 Despite a growing understanding of the processes involved in their generation and their recognized 34 importance on the ocean functioning, the overall impact of the submesoscale dynamics on the 35 vertical tracer fluxes remains ambiguous. It involves competing effects whose balance is likely to 36 be region dependent. On the one hand, the prevailing paradigm is that, at submesoscale, the 37 breakdown of the geostrophic balance leads to large vertical velocities that increase the vertical 38 exchanges between the mixed-layer and the ocean below (e.g. Klein and Lapevre, 2009, Zhong 39 and Bracco 2013, Rosso et al. 2014). On the other hand, the submesoscale eddy tracer fluxes also 40 act to vertically restratify the upper ocean, limiting the deepening of the mixed-layer under surface 41 cooling conditions and speeding up the restratification (Fox-Kemper et al. 2008, Capet et al. 2008, 42 Couvelard et al. 2015). Such control on the mixed-layer depth may in turn limit the vertical tracer 43 exchanges between the subsurface and the mixed-layer. Observations cannot settle this issue but 44 some studies have recently downplayed the role of submesoscale vertical tracer fluxes between the 45 interior ocean and the near-surface (e.g., Ascani et al. 2013). 46 The ubiquity and seasonal variability of small-scale eddies and filaments in the Gulf of Mexico 47 (GoM) is illustrated by the winter and summer surface vorticity snapshots displayed in Figure 1. Despite its potential importance in driving tracers or pollutant dispersion, air-sea exchanges and 48 49 the biogeochemistry of the region, the interest in submesoscale eddy activity in the GoM is quite

- 50 recent and mainly focused on the northern part of the basin (Zhong and Bracco 2013, Luo et al.
- 51 2016, Bracco et al. 2016). Comparing two simulations of the GoM at 5 and 1 km horizontal
- 52 resolution, Zhong and Bracco (2013) found that vertical dispersion of neutrally buoyant water
- 53 parcels increases when increasing the resolution. More recently, Luo et al. (2016) evidenced from
- a high resolution model (1.6 km horizontal resolution) two seasonal peaks of submesoscale activity
 in the northern GoM: a winter peak associated with deep mixed-layers and large available potential
- 56 energy and a summer peak due to the intense lateral density gradients created by the Mississippi
- 57 plume (see also Hetland 2017).

58 The main focus of this study is to clarify the role played by the mesoscale and submesoscale 59 turbulence in controlling the exchanges between the surface and the subsurface of the GoM. Three 60 regional simulations, including a passive tracer sub-model, were performed over a range of 61 horizontal resolution (1/4°, 1/12°, 1/36°) covering eddy present to submesoscale-rich regimes for 62 a period of 20 years. The model experiments are described in Section 2. The mean and seasonal 63 distribution of the meso- to submesoscale activity in the GoM are described from along-track 64 altimetry and our highest resolution simulation (1/36°) in Section 3. The impact of the submesoscale activity on the vertical fluxes of a passive tracer initialized with a vertical profile 65 66 typical of winter nutrient distribution is presented and discussed in Section 4. The results are summarized and discussed in Section 5. 67

68 2. Model and simulations description

69 The numerical code is the oceanic component of the Nucleus for European Modeling of the Ocean 70 program (NEMO3.6, Madec, 2016). It solves the three dimensional primitive equations in 71 spherical coordinates discretized on a C-grid and fixed vertical levels (z-coordinate). Temperature 72 and salinity are advected using a Total Variance Dissipation scheme (TVD) with nearly horizontal 73 diffusion parameterized as a Laplacian isopycnal diffusion. The horizontal diffusion of momentum 74 is implicit by using the third order upstream advection scheme UP3. The vertical diffusion 75 coefficients are given by a Generic Length Scale (GLS) scheme with a k- ε turbulent closure 76 (Umlauf and Burchard 2003, Reffray et al. 2015). Bottom friction is quadratic with a bottom drag coefficient of 10⁻³ and free slip boundary conditions are applied at the lateral boundaries. The 77 78 temporal integration is carried out with a modified Leap Frog Asselin Filter (Leclair and Madec 79 2009), with a coefficient of 0.1. The free-surface is solved using a time-splitting technique with 80 the barotropic part of the dynamical equations integrated explicitly with 60 sub time-steps of the 81 baroclinic time step.

82 The model is forced at its lateral boundaries with daily outputs from the MERCATOR global 83 reanalysis GLORYS2V3. The open boundary conditions radiate perturbations out of the domain 84 and relax the model variables to 1-day averages of the reanalysis data. Details of the method are 85 given in Madec (2016). At the surface, the atmospheric fluxes of momentum, heat and freshwater 86 are computed by bulk formulae (Large and Yeager, 2009). The model is forced with DFS5.2 87 product (Dussin et al. 2016) which is based on ERA-Interim reanalysis and consists of 3-hour 88 fields of wind, atmospheric temperature and humidity, and daily fields of long, short wave 89 radiation and precipitation. The shortwave radiation forcing is modulated by a theoretical diurnal 90 cycle. A monthly climatological runoff based on the data set of Dai and Trenberth (2002) is 91 prescribed near the river mouths as a surface freshwater flux with increased vertical mixing in the

92 upper 10 meters.

93 The regional model configuration consists of a domain encompassing the GoM and the Cayman 94 See (from 0.89W to 789W and from 1.49W to 219W) with 75 levels in the continue (mid-12).

Sea (from 98°W to 78°W and from 14°N to 31°N), with 75 levels in the vertical (with 12 levels in the upper 20 meters and 24 levels in the upper 100 meters). Three numerical experiments at ¹/₄°, 1/12° and 1/36° horizontal resolution have been conducted. Their main characteristics and differences are summarized in Table 1. The three simulations have been integrated from 1 January 1993 to 31 December 2012. They have been initialized using geostrophic balance and the temperature and salinity fields obtained from ¹/₄° outputs of MERCATOR GLORYS2V3

- reanalysis at day 1 January 1993. The 1/36° configuration has been used in Jouanno et al. (2016a)
- 101 to study the formation of short-period Loop Current Frontal Eddies (LCFEs). We refer the reader
- 102 to this publication for a validation of the background state (currents and eddy kinetic energy; EKE)
- and a detailed comparison of the model LCFEs with in-situ and satellite observations.

104 In order to quantify the impact of the high frequency winds on the vertical exchanges of tracers, 105 three parallel simulations at $1/4^{\circ}$, $1/12^{\circ}$ and $1/36^{\circ}$ were carried out with monthly wind stress. In

- 106 this set of simulations (which will be referred as MONTHLY) the surface wind stress is computed
- 107 using monthly averages of the 3-hourly DFS5.2 winds used in the reference simulations. Note that
- 108 in order to identify the role of the mechanical contribution of the high frequency winds (and not
- 109 their direct impact on the air-sea fluxes), the air-sea heat and freshwater fluxes (specifically the
- 110 latent and sensible heat fluxes and the evaporation) in MONTHLY simulations are computed using
- 111 the high-frequency DFS5.2 winds just like in the reference simulations.
- 112 The simulations also include a passive tracer $TR_{nutrient}$ freely transported by the flow and entrained
- by vertical mixing. The passive tracers are advected and diffused with the same scheme employed
- 114 with the active tracers (see above). Following the strategy of Levy et al. (2012), the tracer is
- initially uniformly distributed in the horizontal and, as real nutrients do, increases with depth from 0 at the surface to 1 at 200 m depth and beyond (i.e. below the euphotic depth) in a hyperbolic
- 116 0 at the surface to 1 at 200 m depth and beyond (i.e. below the euphotic depth) in a hyperbolic 117 tangent manner. The tracer is initialized at days Jan 1, Mar 15, May 27, Aug 8, and Oct 20 of each
- 118 year and integrated each time for a period of 73 days. This sums to a total of 100 realizations for
- the period 1993-2012. The connection between the dynamics of this passive tracer and that of
- 119 the period 1995-2012. The connection between the dynamics of this passive tracer and that (
- 120 nutrients in the real ocean is discussed in the conclusions.

3. Meso- and submesoscale surface energy distribution in the Gulf of Mexico: comparison between model and along track altimetry

123 **3.1 Along track sea level data**

124 In order to asses the realism of the meso and submesoscale field resolved by the high resolution 125 GOLFO36 model, we compared surface properties with two different along-track altimeters: 126 Jason2 (for the period 2009-2012) and Saral/Altika (period 2013-2015). We choose these 127 altimeters because of their complete coverage of the area of interest, their performance and their 128 low noise level in the submesoscale range (e.g. Dufau et al. 2016). As in Le Hénaff et al. (2014) 129 we use 1Hz delayed time level-3 sea surface height data not filtered, nor sub-sampled. Jason2 has 130 a repeat period of 10 days for each track, while Altika has a repeat period of 35 days. We computed 131 spectra of both altimetric datasets as follows. Tracks with gaps larger than 18km were removed. 132 The spatial coverage is optimized by segmenting the along track signal into slices of ~ 600 km (as 133 in Dufau et al. 2016) with 50% of authorized overlap between slices. One consequence is that the 134 offshore areas have more weight in the averaged spectra than coastal areas. The model data are 135 interpolated along the tracks with repeat periods similar to the observations.

- 136 The wavenumber sea level spectra from Jason2 and Altika exhibit spectral slopes close to k^{-4}
- 137 (Figure 2a). This slope is consistent with global analysis of Xu and Fu (2012) who found a slope
- 138 close to k^{-4} in the subtropical areas. At wavelengths below 100 km, the model spectral slope is
- 139 steeper than the altimetry slope. In this range, uncertainties are large and there are many possible

140 sources of discrepancy between model and altimeter spectra: i) instrumental noise that increase the

- 141 level of energy seen in altimetry at short scales (particularly so for Jason2, which shows
- 142 significantly more energy than AltiKA there), ii) the sea surface signature of internal tides that are
- 143 not represented in our model configuration and presumably have a significant effect on sea level
- 144 fine-scale variability (Shriver et al. 2012), iii) the use of model daily averages which smooth the
- sea level gradients, and iv) limited effective resolution of our 1/36° grid. At wavelengths between
- 146 600 km and 100 km, the model sea level spectral slope is in very good agreement with the satellite
- 147 observations, suggesting that the model adequately represents KE transfers across scale in this
- 148 wavelength range.

149 **3.2 Kinetic energy**

Along-track kinetic energy has been computed using i) geostrophic velocities obtained from altimetry sea level (Figures 3a,b) ii) geostrophic velocities obtained from model (GOLFO36) sea level (Figures 3c,d), and iii) model (GOLFO36) surface velocities normal to the track (Figures 3e,f). The sea level obtained from the altimeters has been low-pass filtered before applying geostrophy in order to remove the white noise at short wavelength (cutoff wavelength of 40 km). In each case (model and observations), the velocities were high-pass filtered prior to compute the kinetic energy in order to retain only wavelengths below 200 km. Hereinafter, the resulting kinetic

157 energy, representative of the energy below the mesoscale, will be referred as $KE_{\lambda \leq 200 \text{km.}}$

- 158 As expected from previous studies (e.g. Verron et al. 2015), along track $KE_{\lambda < 200 \text{km}}$ in the central 159 GoM illustrates that the noise level is lower in Altika (Figure 3b) compared to Jason2 (Figure 3a). The levels of geostrophic $KE_{\lambda < 200 \text{km}}$ in the model compare well with Altika. Comparison between 160 model based geostrophic KE_{λ <200km} (Figure 3c,d) and KE_{λ <200km} inferred directly from the model 161 162 velocity normal to the track (Figures 3e,f) suggests that the geostrophic approximation allows to 163 catch most of the spatial variability of the small scale features resolved by the model. High levels 164 of $KE_{\lambda < 200 \text{km}}$ are observed in the Loop Current area, as expected by the ubiquitous presence of Loop Current Frontal Eddies (e.g. Jouanno et al. 2016a). The maximum of $KE_{\lambda < 200 \text{km}}$ in the 165 northwestern GoM (near the shelf between 27°N-28°N and 94°W-88°W) described in Luo et al. 166 167 (2016) using a model at 1-km resolution is also reproduced in our 1/36° simulation (Figure 3). It 168 is worth mentioning that the altimetry data does not provide evidences of this maximum (Figure 169 3a,b). The scales involved (O (50) km) may be below the effective resolution of the altimeter and 170 further observations may be required to confirm the presence of enhanced submesoscale activity near the northern GoM shelf. 171
- 172 Several regions with increased short-scale turbulence activity are identified from the temporal

average of model $KE_{\lambda < 200 \text{km}}$ (computed using model surface velocity; Figure 4a). The seasonal

173 average of model $\text{KL}_{\lambda < 200 \text{km}}$ (compared using model surface velocity, Figure 4*a*). The seasonal

- 174 variability of $KE_{\lambda < 200 km}$ present different regimes of variability. In the northwestern GoM, the
- 175 GOLFO36 model shows a semi-annual cycle of $KE_{\lambda < 200 \text{km}}$ (Figure 4b). As mentioned in the
- 176 introduction, such seasonal evolution has been described in Luo et al. (2016) with a higher

177 resolution model (~1.6km). It consists of a dominant winter peak associated with the thermal front 178 at the shelf break and a secondary late summer peak displaced ~100 km northward on the shelf 179 and associated with the haline front of the Mississippi plume (not shown). In the open ocean, both 180 in the western GoM (Figure 4c) and in the Loop Current area (Figure 4d), $KE_{\lambda < 200 \text{km}}$ follows an 181 annual cycle with a peak of energy in winter that follows the mixed-layer depth seasonal cycle 182 with one-month lag (not shown). Such behavior is in agreement with increased small scale 183 turbulence in the deep winter mixed-layer (Sasaki et al. 2014, Callies et al. 2015). The 1-month 184 lag can be interpreted as a consequence of meso/submesoscale energy build-up through available 185 potential energy to eddy kinetic energy conversion by mixed-layer instabilities (which is directly 186 proportional to mixed layer depth; Fox-Kemper et al 2008, Capet et al 2008). In the Loop Current 187 area, the larger density gradients created during winter between the warm waters advected from 188 the Caribbean Sea and the colder waters of the GoM, may also favor higher submesoscale activity 189 in winter. In contrast with the above mentioned regions, the eastern shelf of the Campeche Bank 190 shows a very weak semi-annual cycle, with more energy during winter and summer and a large 191 standard deviation. This may be partly explained by the seasonal cycle of LCFEs dynamics along 192 the Campeche Bank highlighted in Jouanno et al. (2016a) and associated with the seasonality of 193 the cold surge events over the GoM.

194 4. Vertical tracer exchanges

195 With a horizontal grid spacing of ~3 km, the $1/36^{\circ}$ simulation may only marginally be 196 representative of the whole range of submesoscale dynamics in the GoM (Soufflet et al., 2016); 197 nevertheless, we expect to be able to capture a significant fraction of its signature and effects. At 198 lower resolution ($1/12^{\circ}$ and $1/4^{\circ}$), the kinetic energy at $\lambda < 200$ km is well below the energy levels 199 obtained in the $1/36^{\circ}$, at all scales and in all seasons (Figures 2b and 4). In this section, we will 200 show that this has profound consequences on the vertical exchanges of the tracer.

Average surface tracer concentrations for the three model resolutions are shown after 30 days of integration starting from January 1 of each year from 1993 to 2012 (Figures 5j-1). We first focus on the winter period since this is the season with maximum short-scale turbulence activity (see

204 Section 3.3; Figure 4) and as we will see below this is also the season with most sensitivity of

- 205 mixed-layer depth and vertical tracer exchanges to model resolution. The strongest differences are
- 206 in the northwestern part of the GoM, with larger tracer concentrations with decreasing resolution
- 207 (Figures 5j-l), the difference being much larger between the $1/4^{\circ}$ and $1/12^{\circ}$ simulations than
- between the $1/12^{\circ}$ and $1/36^{\circ}$ simulations (Figure 5 and 6a). These differences in surface tracer
- 209 concentrations appear linked with mixed-layer sensitivity to model resolution: the averaged mixed-
- 210 layer depth in the northwestern part of the GoM is ~15 m larger in the lower resolution run
- 211 compared to the higher resolution run.

- 212 This sensitivity of the mixed-layer depth is related to the presence of small-scale processes that
- actively contribute to the restratification of the upper ocean in the higher resolution simulations.
- As mentioned before, the largest impact of increasing resolution is seen in the northwestern part
- 215 of the GoM, where the deepest mixed layers are found (see Figure 5).

216 The vertical tracer injection in the upper mixed layer seen after 30 days of model integration in the northwestern GoM (area between 24°N-28°N and 96°W-90°W) has a marked seasonal cycle 217 218 (Figure 6), with more intense tracer exchange during winter when the mixed-layer is deep (O (100) 219 m) and vanishing tracer exchange in summer when the mixed-layer is shallow (O (10) m). The 220 sensitivity to model resolution is also greater in winter: increasing model resolution shallows the 221 mixed-layer depth by about 30 meters and limits the tracer exchanges (Figure 6a). The monthly 222 evolution of $KE_{\lambda < 200 \text{km}}$ and mixed-layer depth in the northwestern GoM is given in Figure 7a,b. From January to March, $KE_{\lambda < 200 km}$ shows an annual peak which follows with one-month lag the 223 224 annual peak of mixed-layer depth. The vertical buoyancy flux averaged in the upper 200 meters ($<\overline{w'b'}>$ with w' and b' the vertical velocity and buoyancy anomalies with respect to horizontally 225 226 averaged daily values over the area 24°N-28°N and 96°W-90°W, <·> the spatial averaging over 227 this area, and $\overline{\cdot}$ the temporal averaging; Figure 7a) is also maximum from January to March. The 228 eddy driven upward transport of buoyancy turns the upper layers more buoyant and acts to stratify 229 the upper ocean (e.g., see Fox-Kemper et al. 2008). The vertical buoyancy flux is larger in the 230 1/36° simulation, so eddy-driven restratification is particularly active in this simulation, which 231 reduces the mixed-layer depth and limits tracer exchanges.

It is important to note that the r.m.s. of the vertical velocity is much larger in the 1/36° simulation throughout the year (including in winter; Figure 7c) so we might have expected larger tracer exchanges between the mixed-layer and the subsurface in the high resolution run. The comparison of the different simulations contradicts this scenario and rather suggests that, to first order, the

- 236 main effect of fine-scale dynamics on tracer distribution is through eddy-driven restratification.
- 237 In order to investigate how restratification acts during winter, GOLFO36 high-frequency outputs
- for winter 2011 are analyzed during a period of mixed-layer deepening (Figure 8). In January 2011,
- the GoM was under the influence of repeated northerly wind events (the Central American cold
- surges or "Nortes"; e.g. Schultz et al., 1998) that typically last 2 to 4 days and occur every 6-10
- days. During these events, the air-sea fluxes act to cool down the ocean (Figure 8b), with a heat loss that reaches 800 W m⁻² during night time (e.g. day 12 in Figure 8b). For each of these events,
- 242 loss that reaches 800 with during hight time (e.g. day 12 in Figure 80). For each of these events,
- the mixed-layer deepened (Figure 8c) and the surface concentration of the tracer nutrient were increased (Figure 8d). This confirms that the vertical flux of tracer is strongly intermittent and, in
- the specific case of the GoM, it is primarily associated with intense atmospheric synoptic events.
- 246 The intermittent and chaotic shedding of Loop Current eddies travelling westward is expected to
- also modulate the evolution of mixed-layer and thermocline depths in the western GoM, so the
- direct comparison of the high frequency response in January 2011 in the different simulations may

not be meaningful. To partially overcome this issue, we identified a total of 96 surface cooling 249 250 events on the period 1993-2012, each reaching daily-averaged net air-sea heat flux larger than 400 W m⁻². The composite time evolution of mixed-layer depth and vertical buoyancy flux is shown 251 in Figure 9. The averaged mixed-layer deepening (Figures 9a,b) in response to the cooling events 252 253 (500 W m⁻² in average; Figure 9d) is of order 15 meters. The mixed-layer anomalies (referenced 254 to a mixed-layer depth averaged between days 3 to 5 prior to the peak of the wind event) illustrate 255 that the differences observed between the simulations form in the two days following the wind 256 event (Figure 9b). We link these differences with the eddy buoyancy flux which is particularly 257 active during the few days encompassing the cooling event (Figure 9c). Eddy buoyancy fluxes are 258 largest at the highest resolution as expected. This contributes to the reduction of the mixed layer deepening between the 1/12° and 1/36° solutions. $\langle w'b' \rangle$ values at the higher resolution are in 259 260 line with previous estimates (Boccaletti et al. 2007, Brannigan et al. 2015, Capet et al. 2016) and 261 are expected to further increase with increasing resolution.

262 5. Sensitivity under low-frequency wind-stress forcing

263 The high-frequency fluctuations of the wind stress during the passage of atmospheric fronts in the 264 Gulf of Mexico force surface-intensified near-inertial waves (NIW; e.g. see Chen et al., 1996 or Jarosz et al. 2007 for observations in the Gulf of Mexico). A large fraction of the NIW energy is 265 266 thought to dissipate close beneath the mixed-layer base, increasing vertical shear at the base of the 267 mixed-layer and therefore causing entrainment (e.g., Price 1981; Cuypers et al. 2013; Jouanno et 268 al, 2016b). In Section 4, we have shown that high-frequency winds shape the mixed-layer depth 269 and contribute to the upward flux of tracer. One question that arises is how much of the mixed-270 layer deepening is simply due to mixed-layer convection in response to large surface cooling (due 271 to air-sea heat fluxes) and how much is due to the mixing induced by the dissipation of the NIW 272 at the mixed layer base. The presence of small scale eddy vorticity structures is known to 273 accelerates the downward propagation of the wind-forced near-inertial waves (Danioux et al. 274 2011), so we expect some sensitivity of the NIW induced mixing to the model horizontal 275 resolution.

276 This is investigated with the second set of simulations, forced with monthly wind stress, to remove 277 inputs by the winds at the near-inertial scale (referred to as MONTHLY - see Section 2 for details). In these simulations there is no more wind driven input of near-inertial energy. As shown in Figure 278 279 10a,b, the winter mixed-layer depths are slightly reduced when removing the high-frequency 280 winds (reduction of MLD of 5% at $1/4^\circ$, 13% at $1/12^\circ$, and 12% at $1/36^\circ$). The r.m.s. of the vertical 281 velocity (Figures 10c,d) and the vertical flux of buoyancy (Figures 10e,f) in the upper 100 meters 282 are also reduced at all the resolutions with larger differences found, again, for the 1/36° 283 simulations. These results are consistent with those of Jouanno et al (2016b): for near-inertial 284 processes to influence the ocean subsurface it is necessary to, at least, properly resolve the

285 mesoscale activity, which in turn permits the downward propagation of near-inertial motions. Once 286 mesoscale activity is resolved, near-inertial dissipation in the upper thermocline can have a 287 significant impact. We note for example that there is more passive tracer depletion in the subsurface at $1/36^{\circ}$ and high frequency forcings than at $1/12^{\circ}$ with monthly winds. Likewise, we 288 289 also note that the reduction of tracer fluxes exchanges between $1/4^{\circ}$ and $1/36^{\circ}$ resolution is 290 significantly less with high frequency forcings. These comparisons illustrate the complexity of 291 submesoscale impacts on vertical fluxes of properties, with antagonistic processes that respond 292 differently to resolution increase and background conditions (e.g., during summer, removing the 293 high frequency wind stress has almost no impact on the mixed-layer depth or on the upward flux 294 of tracer, not shown).

295 6. Conclusions

This work investigates the role played by the meso- and submesoscale turbulence on the exchanges of passive tracers in the upper layers of the GoM using a hierarchy of regional simulations with horizontal resolution of 1/4°, 1/12° and 1/36° covering eddy present to submesoscale-rich regimes.

299 A careful evaluation of the model mesoscale activity in the GoM has been carried out using along-

300 track data from recent altimeters. At scales larger than 100km, comparison of our higher resolution

301 run with along-track altimetry shows that the model adequately represents the mesoscale field and

302 spectral slope. But such exercise comes up against the high noise levels contained in altimetry

303 products at scales below 50km, which did not allow us to evaluate the realism of the submesoscale

304 field reproduced by the higher resolution simulation.

305 In agreement with the present understanding of submesoscale dynamics, vertical velocity variance 306 in the northwestern GoM increases strongly for increasing horizontal resolution (e.g. Zhong and 307 Braco, 2013). On the other hand, contrary to many previous studies demonstrating that 308 submesoscales enhance surface-subsurface exchanges (through frontogenetic circulations; Klein 309 and Lapeyre 2009, Lévy et al. 2001, Spall and Richards 2000, Mahadevan and Archer 2000) we 310 find that explicitly resolving submesoscale processes may actually limit the exchanges between 311 the surface and the subsurface. Precisely, as shown in our study, submesoscale impacts vertical exchanges of properties in several antagonistic ways. At dx=1km, fluxes of buoyancy are enhanced 312 313 in the upper ocean (compared to lower resolution simulations), which yields a stronger upper ocean 314 stratification (Couvelard et al, 2015), hence a reduction of vertical fluxes of passive or 315 biogeochemical tracers. On the other hand, enhanced resolution favors the penetration of near-316 inertial motions below the mixed layer which tends to strengthen mixing and vertical exchanges. 317 The outcome of these two competing effects depends on the particular distribution of the 318 considered tracer, on the amount of near-inertial energy deposited in the mixed layer, on the Brunt-319 Vaisala profile, etc. With numerical resolution around a few kilometers our passive tracer 320 experiments suggest that both effects are of comparable magnitude.

- This seems to go against the current view that submesoscales would increase nutrient fluxes into the euphotic layer (Lévy et al. 2001, Spall and Richards 2000, Mahadevan and Archer 2000). However, interpretation of our results in terms of real biogeochemical tracers (typically nutrients) should be done with caution. In an equilibrated simulation with explicit representation of biogeochemical sources and sinks the sensitivity to resolution may be felt somewhat differently
- than in our simulations. Over time, the distribution of nutrients may indeed adjust to stratification,
- 327 and thus, at high resolution, develop larger values in the depth range where fluxes are reduced
- 328 compared to lower resolution. On the other hand, we expect the winter-time vertical distribution
- of nutrients in the GoM to be also controlled by external processes, *e.g.*, subsurface replenishment
- through lateral transport by the mean flow and mesoscale turbulence in summer and fall. In such
- 331 conditions, our numerical protocol where a unique initial tracer distribution is prescribed with no
- dependence to resolution might thus provide useful insight into submesoscale impact on nutrient
- 333 fluxes. Analyses of coupled physics-biogeochemistry simulations will be needed to confirm that
- GoM winter submesoscales are detrimental to the injection of subsurface nutrients.

Acknowledgements: This study was supported by CIGOM. We acknowledge the provision of supercomputing facilities by CICESE and GENCI project GEN7298. Altimetry data were produced by Salto/Duacs and distributed by AVISO, with support from CNES. We are grateful to Rachid Benshila and Jerome Chanut for their help with the ocean model. The along-track altimetry data were pre-processed with the Python package py-altimetry-0.3.2.

340 References

- Ascani, F., K.J. Richards, E. Firing, S. Grant, K.S. Johnson, Y. Jia, R. Lukas and D.M. Karl
 (2013). Physical and biological controls of nitrate concentrations in the upper subtropical
 North Pacific Ocean. Deep Sea Research Part II: Topical Studies in Oceanography, 93, 119134.
- Boccaletti, G., R. Ferrari, and B. Fox-Kemper (2007). Mixed layer instabilities and restratification.
 J. Phys. Oceanogr., 2228–2250.
- Bracco, A., J. Choi, K. Joshi, H. Luo, and J.C. McWilliams (2016). Submesoscale currents in the
 Northern Gulf of Mexico: deep phenomena and dispersion over the continental slope. Ocean
 Modelling, 101, 43-58.
- Brannigan, L., D. P. Marshall, A. Naveira-Garabato, and A. G. Nurser (2015). The seasonal cycle
 of submesoscale flows. Ocean Modelling, 92, 69-84.
- Callies, J., R. Ferrari, J.M. Klymak, and J. Gula (2015). Seasonality in submesoscale turbulence.
 Nature communications, 6.

- Capet, X., J.C. McWilliams, M.J. Molemaker, and A.F. Shchepetkin (2008). Mesoscale to sub mesoscale transition in the California current system. Part III: Energy balance and flux. J.
 Phys. Oceanogr. 38, 2256–2269.
- Capet, X., G. Roullet, P. Klein, and G. Maze (2016). Intensification of upper ocean submesoscale
 turbulence through Charney baroclinic instability. J. Phys. Oceanog., 46, 3365–3384, doi:
 10.1175/JPO-D-16-0050.1.
- Chen, C., R. O. Reid, and W. Nowlin, (1996). Near-inertial oscillations over the Texas-Louisiana
 shelf. J. Geophys. Res., 101(C2): 3509–3524.
- Couvelard, X., F. Dumas, V. Garnier, A. L. Ponte, C. Talandier, and A. M. Treguier (2015). Mixed
 Layer formation and restratification in presence of mesoscale and submesoscale turbulence.
 Ocean Modelling, 96, 243-253.
- Dai, A. and K. E. Trenberth (2002). Estimates of freshwater discharge from continents: Latitudinal
 and seasonal variations. J. Hydrometeorol., 3, 660-687.

Danioux, E., P. Klein, M. W. Hecht, N. Komori, G. Roullet, and S. Le Gentil (2011). Emergence
 of wind-driven near-inertial waves in the deep ocean triggered by small-scale eddy vorticity
 structures. Journal of Physical Oceanography, 41(7), 1297-1307.

- Dufau, C., M. Orsztynowicz, G. Dibarboure, R. Morrow, and P. Y. Le Traon (2016). Mesoscale
 resolution capability of altimetry: Present and future. Journal of Geophysical Research:
 Oceans, 121(7), 4910-4927.
- Dussin, R., B. Barnier and L. Brodeau (2016). The making of Drakkar forcing set
 DFS5. DRAKKAR/MyOcean Report 01-04-16, LGGE, Grenoble, France.
- Fox-Kemper, B., R. Ferrari, and R. Hallberg (2008). Parameterization of mixed layer eddies. Part
 I: Theory and diagnosis. Journal of Physical Oceanography, 38(6), 1145-1165.
- Hetland, R. D. (2017). Suppression of baroclinic instabilities in buoyancy driven flow over sloping
 bathymetry. Journal of Physical Oceanography, doi:10.1175/JPO-D-15-0240.1.
- Jarosz, E., Z. Hallock, and W. Teague, (2007). Near-inertial currents in the DeSoto Canyon
 Region. Continental Shelf Research, 27, 2407-2426.
- Jouanno, J., X. Capet, G. Madec, G. Roullet, and P. Klein (2016b). Dissipation of the energy
 imparted by mid-latitude storms in the Southern Ocean. Ocean Sci. 12, 743-769,
 doi:10.5194/os-12-743-2016.
- Jouanno, J., J. Ochoa, E. Pallàs-Sanz, J. Sheinbaum, F. Andrade, J. Candela, J.M. Molines (2016a).
 Loop Current Frontal Eddies: formation along the Campeche Bank and impact of coastally
 trapped waves. Journal of Physical Oceanography, 46(11), 3339-3363.

11

- Klein, P. and G. Lapeyre (2009). The oceanic vertical pump induced by mesoscale and
 submesoscale turbulence. Annual Review of Marine Science, 1, 351-375.
- Large, W. G. and S. Yeager (2009). The global climatology of an interannually varying air-sea ux
 data set. Climate Dynamics, 33, 341-364, doi:10.1007/s00382-008-0441-3.
- Leclair, M. and G. Madec (2009). A conservative leapfrog time stepping method. Ocean
 Modelling, 30, 88-94.
- Le Hénaff, M., V. H. Kourafalou, R. Dussurget, and R. Lumpkin (2014). Cyclonic activity in the
 eastern Gulf of Mexico: Characterization from along-track altimetry and in situ drifter
 trajectories. Progress in Oceanography, 120, 120-138.
- Lévy, M., D. Iovino, L. Resplandy, P. Klein, G. Madec, A. M. Tréguier, S. Masson and K.
 Takahashi (2012). Large-scale impacts of submesoscale dynamics on phytoplankton: Local
 and remote effects. Ocean Modelling, 43, 77-93.
- Lévy, M., P. Klein, A. M. Tréguier, D. Iovino, G. Madec, S. Masson, and K. Takahashi (2010).
 Modifications of gyre circulation by sub-mesoscale physics. Ocean Modelling, 34(1), 1-15.
- Lévy M, P. Klein, and A.M. Treguier (2001). Impacts of sub-mesoscale physics on production and
 subduction of phytoplankton in an oligotrophic regime. J. Mar. Res. 59, 535–65.
- Luo, H., A. Bracco, Y. Cardona, and J.C. McWilliams (2016). Submesoscale circulation in the
 northern Gulf of Mexico: surface processes and the impact of the freshwater river input. Ocean
 Modelling, 101, 68-82.
- 406 Madec, G. and the NEMO team (2016). "NEMO ocean engine". Note du Pôle de modélisation,
 407 Institut Pierre-Simon Laplace (IPSL), Paris, France, No 27 ISSN No 1288-1619.
- Mahadevan, A. and D. Archer (2000). Modeling the impact of fronts and mesoscale circulation on
 the nutrient supply and biogeochemistry of the upper ocean. J. Geophys. Res. Oceans
 105:1209–25.
- McWilliams, J. C. (2016). Submesoscale currents in the ocean. In Proc. R. Soc. A, Vol. 472, No.
 2189, p. 20160117.
- Reffray G., R. Bourdalle-Badie and C. Calone (2015). Modelling turbulent vertical mixing
 sensitivity using a 1-D version of NEMO. Geosci. Model Dev., 8, 69–86, doi:10.5194/gmd-869-2015.
- Rosso, I., A. M. Hogg, P. G. Strutton, A. E. Kiss, R. Matear, A. Klocker, and E. van Sebille (2014).
 Vertical transport in the ocean due to sub-mesoscale structures: Impacts in the Kerguelen
 region. Ocean Modelling, 80, 10-23.

- 13
- Sasaki, H., P. Klein, B. Qiu, and Y. Sasai (2014). Impact of oceanic-scale interactions on the
 seasonal modulation of ocean dynamics by the atmosphere. Nat. Commun., 5, 5636,
 doi:10.1038/ncomms6636.
- Schultz, D. M., W. E. Bracken, and L. F. Bosart (1998). Planetary- and synoptic-scale signatures
 ssociated with central american cold surges. Mon. Wea. Rev., 126, 5–27.
- Shriver, J. F., B. K. Arbic, J. G. Richman, R. D. Ray, E. J. Metzger, A. J. Wallcraft, and P. G.
 Timko (2012). An evaluation of the barotropic and internal tides in a high-resolution global
 ocean circulation model. J. Geophys. Res., 117, C10024, doi:10.1029/2012JC008170.
- Soufflet, Y., P. Marchesiello, F. Lemarié, J. Jouanno, X. Capet, L. Debreu, and R. Benshila (2016).
 On effective resolution in ocean models. Ocean Modelling, 98, 36-50.
- Spall S.A. and K.J. Richards (2000). A numerical model of mesoscale frontal instabilities and
 plankton dynamics—I. Model formulation and initial experiments. Deep-Sea Res. I 47:1261–
 301.
- 432 Umlauf, L. and H. Burchard (2003). A generic length-scale equation for geophysical turbulence
 433 models. Journal of Marine Research, 61 (31), 235–265.
- Verron, J., P. Sengenes, J. Lambin, J. Noubel, N. Steunou, A. Guillot, N. Picot, S. Coutin-Faye,
 R. Sharma, R. M. Gairola, D. V. A. R. Murthy, J. G. Richman, D. Griffin, A. Pascual, F. Rémy,
 and P. K. Gupta (2015). The SARAL/AltiKa altimetry satellite mission. Marine Geodesy,
 38(sup1), 2-21.
- Xu, Y. and L. L. Fu (2012). The effects of altimeter instrument noise on the estimation of the
 wavenumber spectrum of sea surface height. Journal of Physical Oceanography, 42(12), 22292233
- 441 Zhong, Y. and A. Bracco (2013). Submesoscale impacts on horizontal and vertical transport in the
- 442 Gulf of Mexico. Journal of Geophysical Research: Oceans, 118(10), 5651-5668.

Name	Horizontal resolution	Δx, Δy	Δt	Laplacian diffusivity coefficient
GOLFO36	1/36°	~2.8 km	150s	$45 \text{ m}^2 \text{ s}^{-1}$
GOLFO12	1/12°	~8.4 km	600s	$135 \text{ m}^2 \text{ s}^{-1}$
GOLFO04	1∕4 °	~25.2 km	1800s	$405 \text{ m}^2 \text{ s}^{-1}$

Table 1: Description of the three reference simulations used in this study.



444 **Figure 1**. Snapshots of surface vorticity simulated by GOLFO36 at days 15 February 2008 (a) and

- 445 15 August 2008 (b). We identify mesoscale structures (the Loop Current, anticyclones close that
- 446 have been shed by the Loop Current), and submesoscale structures including filaments.



447 Figure 2. a) Along track wavenumber spectra from altimetry SSH (Jason2 in black, Altika in red) and GOLFO36 model SSH (along Jason2 tracks in light black and along Altika tracks in light red). 448 449 A multitaper method with taper of 600km has been used. The model data have been interpolated along the satellite tracks. b) KE spectra for the horizontal velocity \mathbf{u}_h at 10-m depth for the three 450 simulations at $1/36^{\circ}$, $1/12^{\circ}$ and $\frac{1}{4}^{\circ}$ plotted as a function of horizontal wavenumber magnitude, k =451 $|\mathbf{k}_{h}|$; so the 2D spectrum is azimuthally integrated in k shells following the method described in 452 Capet al. (2008c). The corresponding variance conserving spectra are shown in the inset (units 10⁻ 453 ⁶ m s⁻²). Daily averages model data from 98W to 82W and 18N to 30N and from 2008 to 2012 454 have been used and the spectra have been temporally averaged. The k^{-2} , k^{-3} and $k^{-5/3}$ spectral slope 455 456 are also shown.



Figure 3. KE [log10 m²s⁻²] in the 40-200km band obtained from along track altimetry SSH (top),
GOLFO36 model SSH (middle), and surface velocity (bottom). The altimetry SSH has been
filtered with a 40km low pass filter before applying geostrophy. The left column is for Jason2 and
the right column is for AltiKa data.



Figure 4. KE (m² s⁻²) computed from high-pass filtered ($\lambda < 200$ km) surface velocity anomalies : mean distribution in GOLFO36 (a) and seasonal cycle of KE_{λ -200km} spatially averaged over four different regions (b-e), for the three simulations (GOLFO36 in red, GOLFO12 in blue and GOLFO04 in black). Daily data from 2003 to 2012 have been used and the enveloppes indicate the monthly s.t.d. of the daily values of KE.



Figure 5. Winter (a) mean surface kinetic energy (m² s⁻²), (b) mean KE_{x-200km} (log10 m² s⁻²) 467

W 86°W 84°W

- computed from high-pass filtered ($\lambda < 200$ km) surface velocity, (c) mean mixed-layer depth 468
- 469 (MLD; unit m), and (d) 5-day mean averaged surface distribution of passive tracers
- "NUTRIENT" after 30 days of transport in the 1/4° (left), 1/12° (center) and 1/36° (right) 470
- 471 simulations. The MLD criterion used in this study corresponds to the depth at which density is
- 0.03 kg m⁻³ larger than the density at 10 m depth. Data from 1 January to 14 March of each year 472
- from 1993 to 2012 are considered and the tracers were initialized at five different dates 473
- 474 (regularly spaced; see text) for each year from 1993 to 2012.



Figure 6. Results of on-line transport experiments of the passive tracer "NUTRIENT" performed for the 1/4°, 1/12° and 1/36° simulations and averaged over the northwestern GoM (96°W-90°W, 24°N-28°N). The tracer was initialized at the five same dates of each year (a-e) from 1993 to 2012. The initial state is uniform in the horizontal, with a vertical gradient on the vertical, increasing from 0 at the surface to 1 at 200 m and below. Vertical profiles of the tracer anomaly (i.e. the difference between the final profile and the initial profile) are shown after 30 days of transport. The horizontal dashed lines indicate the averaged mixed-layer depth for each 30-day period.



Figure 7. Seasonal cycle of (a) surface $\text{EKE}_{s=200\text{km}}$ (N m⁻²) computed from high pass filtered ($\lambda < 200\text{km}$) horizontal velocity anomalies, (b) mixed-layer depth (m), (c) the r.m.s. of the vertical velocity averaged in the upper 200 meters (m s⁻¹) and (d) the eddy buoyancy flux $\langle \overline{w'b'} \rangle$ averaged in the upper 200 meters (m² s⁻³). The seasonal cycles are built using daily data of the three different simulations (GOLFO36 in red, GOLFO12 in blue and GOLFO04 in black) from 1993 to 2012, and were averaged in the northwestern GoM (96°W-90°W, 24°N-28°N). The monthly standard deviation is displayed in light colors.



489 Figure 8. Hourly time series of 1/36° model fields averaged in the northwestern GoM (96°W-

490 90°W, 24°N-28°N) from days 1 to 31 January 2011 : (a) wind stress (N m⁻²), (b) net air-sea

491 fluxes (W m⁻²), (c) mixed-layer depth (m) and (d) surface concentration of tracer "Nutrient"

492 initialized at day 0 in GOLFO36.



Figure 9. Lagged composites based on 96 events identified with upward daily averaged net –airsea heat flux larger than 400 W m⁻², obtained for the reference set of simulations: a) mixed-layer depth (m), b) mixed-layer depth anomaly (m) referenced to the mixed-layer depth averaged between lags -5 to -3 days, c) vertical buoyancy flux $\langle \overline{w'b'} \rangle$ averaged between 0 and 200 meters (m² s⁻³) and d) net air-sea heat flux (W m⁻²; negative means upward). Daily data from 1994 to 2012 have been used.



499 Figure 10. Comparison between the reference set of simulations forced with 3-hours wind stress
500 (REF) and the simulations forced with monthly wind stress (MONTHLY) during winter: a,b)

501 passive tracer anomaly after 30 days of transport during the winter period (as in Figure 6a), c,d)

502 r.m.s. of the vertical velocity (m s⁻¹) and e,f) vertical buoyancy flux $< \overline{w'b'} > (m^2 s^{-3})$. All the

503 diagnostics have been made using data from January 1 to March 14 for year from 1994 to 2012

and were averaged in the northwestern GoM ($96^{\circ}W-90^{\circ}W$, $24^{\circ}N-28^{\circ}N$).