

1 **Mechanisms of heat transport across the Southern**
2 **Greenland continental shelf in two eddy-active**
3 **ocean/sea-ice simulations**

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8 **Key Points:**

- 9 • Cross-shelf heat transport is strongest over the southeast continental shelf in both
10 ocean models
- 11 • Shelf currents bring heat to the southwest shelf, where winds drive off-shelf heat
12 transport
- 13 • Topographic Rossby waves generated near the Denmark Strait impact cross-isobath
14 heat transport

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Abstract

The increased presence of warm Atlantic water on the Greenland continental shelf has been connected to the accelerated melting of the Greenland Ice Sheet, particularly in the southwest and southeast shelf regions. Results from two eddy-permitting coupled ocean-sea ice simulations are used to understand the transport of heat on and off the southern Greenland shelf. The analysis reveals that the region of greatest heat transport onto the shelf is southeast Greenland. On the southwestern shelf, heat is mainly exported from the shelf to the interior basins. Full heat budgets for a series of control volumes on the Greenland shelf are analyzed to identify the mechanisms that drive cross-shelf heat transport. Two mechanisms of shelf-basin heat exchange are explored: wind-driven exchange and topographic Rossby waves generated near the Denmark Strait. The heat and volume transports do not depend directly on the winds, and differences in wind forcing between the two simulations may contribute only slightly to differences in properties on the shelf and in the transport of heat across the shelf. Topographic Rossby waves are observed in both simulations along the southeast shelf break; in both models they modify cross-shelf transport as they propagate clockwise along the shelf. On the southern shelf, warm water is spread to the southeast and then southwest Greenland by coastal and boundary currents.

Plain Language Summary

Melting of the Greenland Ice Sheet has been accelerating in recent decades because of rising ocean and air temperatures. Warm ocean water in the deep basin from the subtropical North Atlantic is separated from the ice sheet margin (glacier termini in the Greenland fjords) by the shallower continental shelf region. In this study we compare two simulations of the ocean and sea ice that represent the currents and eddying motions around Greenland realistically. We identify how and where heat is moved on and off the southern Greenland shelf and consider the results to be robust when they are common to both simulations. We find that warm water mainly moves onto the southeast shelf and off the southwest shelf; the currents on the shelf transport the warm water around the southern tip of Greenland. Winds do not directly move heat on or off the shelf. Near the Denmark strait we identify oscillations, known as topographic Rossby waves, that propagate southwards along the shelf producing intrusions of warm water onto the shelf during their

46 passage. Understanding how warm water reaches the shelf allows us to better understand
47 how the ocean contributes to the melting of the Greenland Ice Sheet.

48 **1 Introduction**

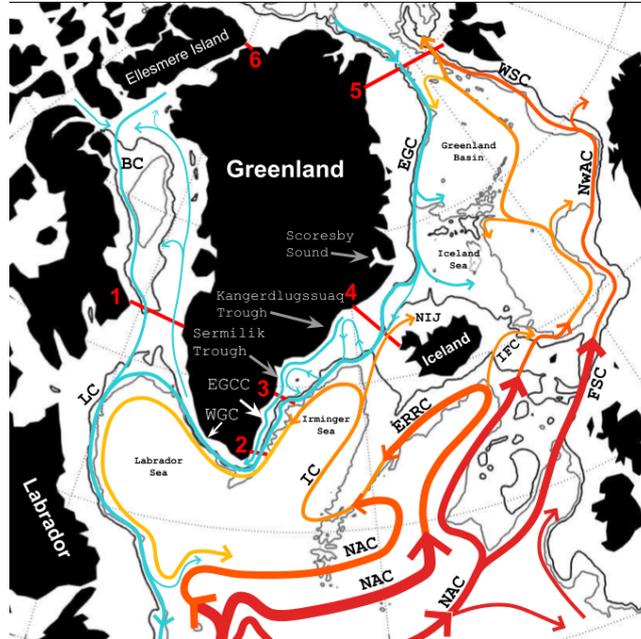
49 Mass loss from both the Greenland and Antarctic Ice Sheets remains the major source
50 of uncertainty in projecting global sea level rise (Meehl et al., 2007). The Greenland Ice
51 Sheet (GIS) is losing mass at an increasing rate, from 51 ± 17 Gt/y in the 1980s to 286
52 ± 20 Gt/y in the 2010s (Mouginot et al., 2019). Since 1972, this mass loss has contributed
53 to 13.7 ± 1.1 mm of global sea level rise (Mouginot et al., 2019). Recently B. Smith et
54 al. (2020) reported a total mass loss of 200 ± 12 GT/y from 2003 to 2019. Projections
55 of sea level rise due to ice sheet mass loss emphasize the short-term (next 100 years) im-
56 portance of the GIS as oceanic and atmospheric temperatures rise (Meehl et al., 2007).
57 The lack of representation of both ice sheet dynamics and connections to the ocean and
58 atmosphere in climate models contributes significantly to the uncertainty of those pro-
59 jections. An estimated 15–25% of total mass loss from the GIS is from melting tide-
60 water glaciers with an additional 15–25% from calving fluxes (Benn et al., 2017); both
61 processes increase as the ocean warms.

62 The margin of the GIS is comprised of both land terminating and marine termi-
63 nating glaciers; the tidewater terminating glaciers are the primary connection between
64 the ocean and the GIS through deep narrow fjords. Warm salty water of Subtropical North
65 Atlantic origin is thought to provide the source of heat needed for ocean-driven melt-
66 ing (Straneo & Heimbach, 2013; Rignot et al., 2012); the co-location of Atlantic-originated
67 water and tidewater glaciers makes the southeastern portion of the GIS particularly vul-
68 nerable to ocean-driven melting (Millan et al., 2018). Over the southeast portion of the
69 GIS, the observed mass loss (Luthcke et al., 2006; van den Broeke et al., 2009; Wouters
70 et al., 2008) is, in part, attributed to warming ocean conditions (Howat et al., 2008), but
71 it is difficult to separate these effects from those of atmospheric warming (Straneo et al.,
72 2013; Hanna et al., 2013). The presence of warm water on the southwest shelf has also
73 been observed (Sutherland et al., 2013; Straneo et al., 2012). Observations from specific
74 glacial fjords have shown warming of ocean water preceding glacial retreat events (Christoffersen
75 et al., 2012; Holland et al., 2008), implying that in some regions heat from the ocean may
76 be the leading driver of ice sheet mass loss. Within fjords, observations have provided

77 estimates of the penetration of warm water to the front of glaciers (Jackson et al., 2014)
78 given the presence of Atlantic water on the shelf.

79 Comprehensive observations of the subpolar and polar oceans are uniquely chal-
80 lenging to obtain (G. C. Smith et al., 2019). Therefore, the available data are often sparse
81 and irregularly sampled. Ocean models can be used to better understand the ocean cir-
82 culation in these regions, as they provide a complete temporal and spatial record. How-
83 ever, the production of realistic simulations in these regions is challenging. There are lim-
84 ited observations (Morlighem et al., 2017; An et al., 2019) of the topography of the Green-
85 land shelf, particularly the fjord bathymetries. In atmospherically forced simulations of
86 the ocean, the quality of the atmospheric forcing will determine the accuracy of the ocean’s
87 response. In addition, simulations that aim to study shelf and near-shelf processes re-
88 quire high resolution in those regions, making them computationally expensive.

99 The Greenland continental shelf is impacted by the fresh and cold water masses
100 exported from the Arctic Ocean as well as the warm and salty water masses advected
101 from the North Atlantic (Figure 1) (after Holliday et al., 2018). Warm water from the
102 subtropical gyre is advected into the subpolar gyre by the North Atlantic Current (NAC),
103 an extension of the Gulf Stream. The NAC consists of multiple northward branches; east-
104 ward branches enter the Nordic seas, while those to the west retrofect to enter the Irminger
105 Current (Holliday et al., 2018). Just south of the Denmark Strait, the Irminger Current
106 retroflects, and its primary branch heads southward along the Greenland continental shelf
107 break. On the Greenland Shelf, from the Fram Strait to Cape Farewell, the East Green-
108 land Current (EGC) flows southward, advecting cold fresh water from the Arctic and
109 seasonal sea ice melt. At Kangerdlugssuaq Trough, the smaller East Greenland Coastal
110 Current (EGCC) develops (Sutherland & Pickart, 2008) onshore of the EGC. The cir-
111 culation along the East Greenland Shelf is characterized by cold fresh shelf waters and
112 by warmer saltier waters off the shelf in the Irminger Basin. Warm water in the basin
113 cannot easily cross onto the shelf because the shelf break is a barrier to onshore trans-
114 port. In an unforced system, a water parcel will conserve potential vorticity by balanc-
115 ing the Coriolis force f and the height of the water column h . To cross the shelf break,
116 changing h by moving from a deep basin to a shallow shelf, a source of potential energy
117 is needed.



89 **Figure 1.** Schematic of circulation in the Subpolar North Atlantic. Major currents are labeled: North Atlantic Current (NAC), East Reykjavik Ridge Current (ERRC), Irminger Current (IC), East Greenland Current (EGC), East Greenland Coastal Current (EGCC), West
 90 Greenland Current (WGC), Labrador Current (LC), Baffin Current (BC), North Icelandic Jet
 91 (NIJ), Iceland-Faroe Current (IFC), Faroe-Shetland Current (FSC), Norwegian Atlantic Current (NwAC), and West Spitsbergen Current (WSC). The major transects used to divide the on-shelf
 92 regions are labeled 1-6: (1) Davis Strait, (2) Cape Farewell Gate, (3) Sermilik Gate, (4) Den-
 93 mark Strait, (5) Fram Strait, and (6) Nares Strait. (after Holliday et al., 2018) with additions
 94 from (Sutherland & Pickart, 2008; Håvik et al., 2017; Saini et al., 2020; Furevik & Nilsen, 2005;
 95 Rossby et al., 2018).

118 In this paper, we aim to understand the processes that could drive cross-isobath
119 exchange at the Greenland Continental Shelf, and we therefore look into previous stud-
120 ies of analogous shelves. For example, the Antarctic Continental Shelf is a similar high-
121 latitude shelf which, in some regions, has cold water on the shelf, strong shelf currents,
122 and warm water in the deep basin in a configuration that is analogous to southeast Green-
123 land. Studies of mechanisms of heat transport onto the Antarctic Continental Shelf have
124 been similarly motivated by the need to better understand ocean-driven melting of the
125 Antarctic Ice Sheet. These studies have explored possible drivers for the on-shelf trans-
126 port of warm Circumpolar Deep Water across the Antarctic continental shelf break. Winds
127 along the Antarctic Slope position the pycnocline at the shelf break to balance the eddy-
128 driven on-shelf transport of Circumpolar Deep Water (Stewart & Thompson, 2015). Weak-
129 ened easterly winds shoal isotherms, while freshening on the shelf opposes the shoaling,
130 re-enforcing the importance of both winds and lateral density gradients in cross-shelf trans-
131 port (Goddard et al., 2017). In regions where the Antarctic Slope Front is strong (from
132 the Weddell Sea to the Ross Sea, including East Antarctica) eddy stirring along the front
133 has been found to be a key mechanism for heat transfer (Stewart et al., 2018). Palóczy,
134 Gille, and McClean (2018) found using the same high-resolution POP simulation used
135 in this study, see Section 2.1 for details, that the mean heat transport onto the Antarc-
136 tic Continental Shelf is controlled by Ekman convergence over the shelf. Coastal troughs
137 play a key role in the pathways of Circumpolar Deep Water onto the continental shelf;
138 Dinniman, Klinck, and Smith Jr (2011) found that warm water intrusions in a trough
139 in the West Antarctic Peninsula were linked to short-duration wind events. The mech-
140 anisms found to transport Circumpolar Deep Water across the Antarctic Slope are likely
141 to be relevant to the Greenland Continental Shelf and provide the motivation for the mech-
142 anisms we will explore to explain cross-shelf heat transport around Greenland.

143 For the Greenland Shelf, thus far there have not been studies akin to the Antarc-
144 tic investigations of mechanisms governing on-shelf heat transport and the resulting warm-
145 ing. A recent study from Gillard, Hu, Myers, Ribergaard, and Lee (2020) focused on troughs
146 around the shelf with a focus on heat reaching fjords. Other model based studies have
147 found high-frequency variability along the southeast Greenland Continental Shelf occurs
148 in high-resolution regional simulations (Moritz et al., 2019; Gelderloos et al., 2021) and
149 cyclonic eddies formed at the Denmark Strait propagate along the shelf break (Moritz
150 et al., 2019). Coastal trapped waves in this region have been identified in a high-resolution

151 simulation from subinertial variability (Gelderloos et al., 2021) though the waves were
152 not specifically defined as topographic Rossby waves.

153 Questions then arise as to how efficiently heat is transported onto the shelf, and
154 in what regions. Understanding which processes drive the transport of warm Atlantic
155 water onto the shelf will help better project what changes to the regional winds, local
156 stratification, or circulation on the shelf could be linked to increased vulnerability of tide-
157 water glaciers to ocean-driven melting. Therefore, our study focuses on understanding
158 both the mechanisms governing the resulting pattern of temperature variability on the
159 shelf using two coupled ocean–sea-ice simulations performed with the Parallel Ocean Pro-
160 gram (POP) and the HYbrid Coordinate Ocean Model (HYCOM), respectively. By com-
161 paring temperature on the shelf and the cross-shelf heat transport in the two simulations,
162 we are able to gain insight into the dominant mechanisms of shelf–basin exchange. The
163 two simulations are configured differently and use different atmospheric forcing and there-
164 fore are independent experiments in which the mechanisms that drive on-shelf heat trans-
165 port and shelf–basin exchange are explored.

166 In Section 2, we begin with a description of the models and, in Section 3, compare
167 model results to observations. Next, in Section 4, we calculate cross-isobath volume and
168 heat transports and then examine the spatial patterns of temperature and cross-shelf heat
169 transport in southern Greenland. We look at the heat flux through key gates around the
170 shelf to gain insights into the heat flux of the East Greenland Current system on the con-
171 tinental shelf, with the understanding that our models resolve limited dynamics on the
172 shelf. We find seasonal and daily variability, motivating the exploration of two mecha-
173 nisms. In Section 5, a primary cross-isobath exchange mechanism is proposed for each
174 timescale: wind-driven exchange on the seasonal scale and topographically trapped waves
175 that modify cross-shelf exchange on a multi-day timescale. Finally, we discuss the ro-
176 bust dynamics identified by comparing these two independent simulations.

177 **2 Model Description**

178 We compare two coupled ocean sea-ice models with horizontal resolutions compa-
179 rable to the first baroclinic Rossby radius of deformation in this region (4–8 km in the
180 deep ocean). Models with this resolution are classified as “eddy-permitting” (Dukhovskoy
181 et al., 2016; Nurser & Bacon, 2014). The effective grid spacing in POP is ~ 5 –6 km and

182 \sim 4-5 km in HYCOM. The first baroclinic Rossby radius is smaller on the continental
 183 shelf (2-4 km Nurser & Bacon, 2014); both models have limited ability to capture smaller
 184 mesoscale processes on the shelf, in the next section a description of what shelf mesoscale
 185 processes are that captured is presented. Each model is forced by a different set of at-
 186 mospheric observations, but neither assimilates data. This allows each model to act as
 187 an independent representation of the dynamics in this region. Both models are coupled
 188 to the same sea-ice model and do not include any representation of freshwater from GIS
 189 melt. Each simulation is described in further detail in the next subsections.

190 **2.1 POP 0.1° - CICE 4**

191 In this study, we use results from a global 62-year (1948-2009) simulation of the
 192 Parallel Ocean Program version 2 (POP2; Dukowicz & Smith, 1994) and the Commu-
 193 nity Ice Code version 4 (CICE4; E. C. Hunke et al., 2010) coupled together in the Com-
 194 munity Earth System Model (CESM; Hurrell et al., 2013) version 1.2 framework (McClean
 195 et al., 2018). The ocean and sea-ice components are coupled every 6 hours using Flux
 196 Coupler 7 (CPL7, Craig et al., 2012). For further details on this simulation, see (Wang
 197 et al., 2018; Palóczy et al., 2018, 2020; Castillo-Trujillo et al., 2021; Arzeno-Soltero et
 198 al., 2021). This simulation is referred to as POP from here on in the text.

199 The ocean and sea-ice models are on a 0.1° tripolar grid with nominal horizontal
 200 resolution of \sim 5–6 km in our study region. POP has 42 non-uniformly spaced vertical
 201 levels; they range from 10-m spacing at the surface to 250 m in the deep ocean. At 800 m
 202 depth, the vertical spacing is approximately 200 m. The bathymetry is based on ETOPO2
 203 with minor modifications in the Arctic (more details are given by McClean et al. (2011)).
 204 Partial bottom cells are used to more smoothly represent the bathymetry. The ocean model
 205 has an implicit free surface and is globally volume conserving.

206 Sub-gridscale horizontal mixing of momentum and tracers in POP is represented
 207 using bi-harmonic operators. The horizontal viscosity and diffusivity coefficients depend
 208 on the horizontal grid spacing, with equatorial values of $-2.7 \times 10^{10} \text{ m}^4 \text{ s}^{-1}$ and $-0.3 \times$
 209 $10^{10} \text{ m}^4 \text{ s}^{-1}$, respectively, varying with the cube of the average grid cell length. The K-
 210 profile parameterization (Large et al., 1994) is used to represent vertical mixing. This
 211 simulation setup does not include any explicit tidal forcing or additional mixing from
 212 tidal dissipation.

213 POP's ocean properties, potential temperature and salinity, were initialized from
214 the World Hydrographic Program Special Analysis Center climatology (Gouretski & Kolter-
215 mann, 2004). A 17-day stand-alone (without CICE4) spin-up integration was run using
216 very short time steps to allow high-amplitude transients to equilibrate. The coupled ocean/sea
217 ice simulation was then initialized from the end state of the 17-day POP simulation. The
218 atmospheric forcing is given by the Coordinated Ocean-ice Reference Experiment-II (CORE-
219 II) corrected interannually varying fluxes (CIAF; Large & Yeager, 2009)) and has a hor-
220 izontal resolution of $\sim 1.9^\circ$. Starting in 1980, climatological values of radiation and pre-
221 cipitation are replaced with time-varying data based on observations. Ocean surface evap-
222 oration and precipitation fluxes and runoff are implemented using virtual salt fluxes; a
223 surface salinity restoring condition with an effective timescale of about 4 years limits model
224 drift.

225 Sea-ice velocities are defined by the dynamic component of CICE4 based on the
226 elastic-viscous-plastic (EVP) rheology of E. Hunke and Dukowicz (1997). Sea-ice and
227 snow growth rates are defined by the thermodynamic component of CICE4 based on the
228 energy conserving sea ice of Bitz and Lipscomb (1999). The Delta-Eddington multiple
229 scattering parameterization is used to represent solar radiation transfer between ice and
230 snow (Briegleb & Light, 2007).

231 The sea-ice state was initialized from a uniform 2-m thick layer with the ice edge
232 defined by the January climatological ice edge from Special Sensor Microwave Imager
233 (SSM/I) observations. In the POP simulation, the Arctic sea ice adjusted over a longer
234 period than in the Antarctic, with potentially excessive melt and export through the Fram
235 Strait during the 1960s. This long adjustment and excessive export could be a transient
236 solution from the initialization of the simulation. Another alternative is that the clima-
237 tological period of the atmospheric forcing (1948-1980) exaggerates a large-scale forc-
238 ing that could increase sea-ice export.

239 We use output from the last five years of the simulation for our analysis. For this
240 period, daily averages of the model variables were saved.

241 **2.2 0.08° Arctic Ocean HYCOM - CICE 4**

242 The second model used in this study results from numerical experiments by Dukhovskoy
243 et al. (2019) conducted using regional 0.08° Arctic Ocean HYbrid Coordinate Ocean Model

244 (HYCOM) (Bleck, 2002; Chassignet et al., 2003, 2007) coupled to CICE4. This simu-
245 lation is referred to as HYCOM in the text.

246 The model domain is a subset of the global HYCOM (Chassignet et al., 2009; Met-
247 zger et al., 2014) north of 38°N . The computational grid of the 0.08° HYCOM-CICE
248 is a Mercator projection from the southern boundary to 47°N . North of 47°N , it employs
249 an orthogonal curvilinear Arctic dipole grid (Murray, 1996). The model has effective spac-
250 ing of $\sim 4\text{--}5$ km in the Subpolar North Atlantic. The model topography is derived from
251 the Naval Research Laboratory Digital Bathymetry Data Base 2-minute resolution (NRL
252 DBDB2). In the current configuration, HYCOM employs a vertical grid with 41 hybrid
253 layers that provide higher resolution in the upper 1500 m. This simulation is one-way
254 nested within the 0.08° Global HYCOM +Navy Coupled Ocean Data Assimilation (NCODA)
255 3.0 reanalysis (Metzger et al., 2014) (for 1993–2005) and Global Ocean Forecasting Sys-
256 tem (GOFS) 3.1 analysis (for 2006–2016).

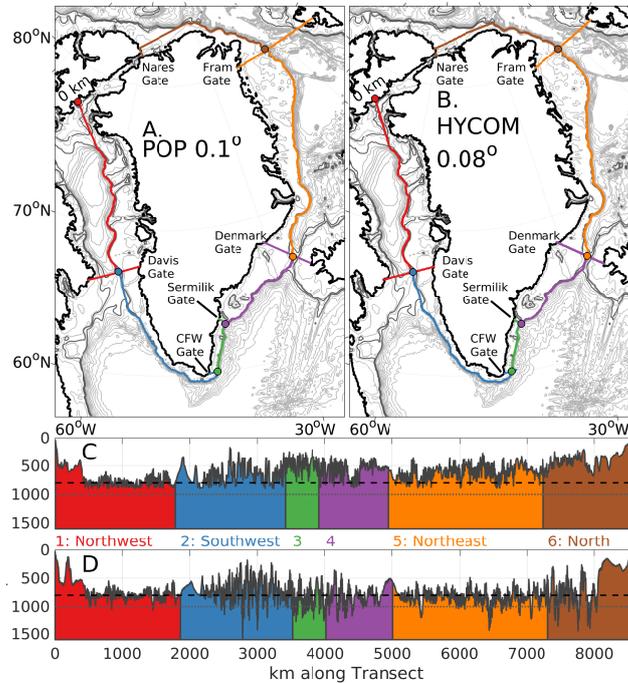
257 Atmospheric forcing fields are obtained from the National Centers for Environmen-
258 tal Prediction (NCEP) Climate Forecast System Reanalysis (CFSR) (Saha et al., 2010)
259 for 1993–2011 and CFSv2 (Saha et al., 2014) for 2012–2016. More details on the model
260 configuration and computational grid as well as model validation and analysis of the model
261 experiments are given by Dukhovskoy et al. (2019, 2021).

262 **3 Model Validation**

263 **3.1 Continental Shelf Control Volumes**

264 To understand how warm Atlantic water crosses onto the shelf and where it is present,
265 the shelf and shelf break must be clearly defined. Shallow straits and deep troughs make
266 choosing a single isobath for the shelf break challenging. However, we found the 800 m
267 isobath to be representative of the shelf break; our results were not sensitive to small changes
268 in the choice of isobath. The contour surrounds Greenland, starting and ending at the
269 Nares Strait connecting to Ellesmere Island (see Figure 2). The exact depths of the shelf
270 break in each model (see Figure 2) show how the bathymetry of the two simulations dif-
271 fers. See supplemental material for a detailed map of the Southeast region highlighting
272 the troughs and small scale bathymetry.

273 In addition to the shelf break, we define six control volumes to examine spatial dif-
274 ferences in cross-isobath transports and properties on the shelf. The boundaries are spec-
275 ified in the supplemental material.



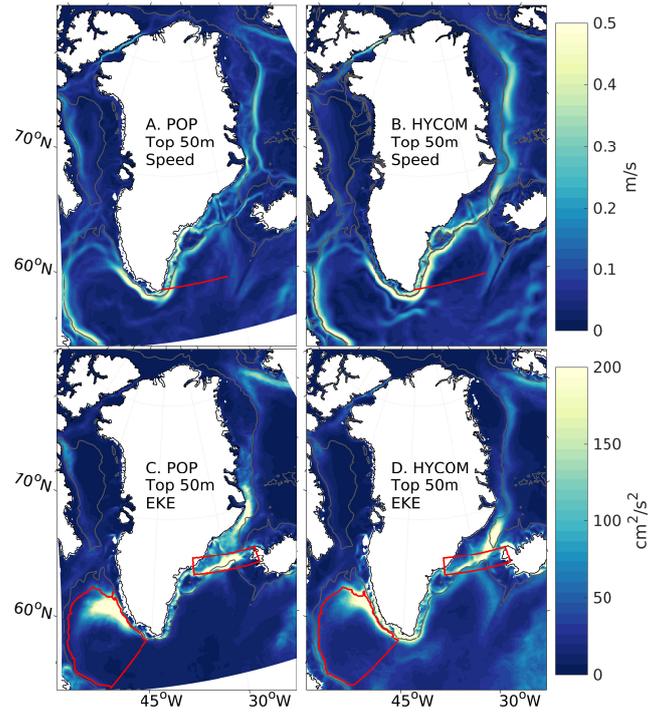
264 **Figure 2.** Map of circum-Greenland transects in (A) POP and (B) HYCOM, subdivided at
 265 the major straits and gates and plotted over the regions' bathymetries. The 800 m isobath is
 266 followed in both models, with the exact depth along the transect plotted for (C) POP and (D)
 267 HYCOM with the regions numbered; the black dashed line is 800m and the gray dashed line
 268 is 1000m. Shelf regions are: (1) Northwest, (2) Southwest, (3) Narrow Shelf, (4) Wide Shelf,
 269 (5) Northeast, and (6) North. Together Sections 3 and 4 constitute the Southeast region. The
 270 color of the contour in each region corresponds to the bathymetry plotted for that region. A
 271 regional map of the Southeast region directly comparing the two bathymetries is provided in the
 272 supplemental materials.

284 ified at major straits or “gates” to differentiate heat transport regionally. The fluxes at
 285 the gates are also computed, providing a representation of the circulation on the shelf.
 286 However, neither model fully resolves the sub-mesoscale dynamics that are important
 287 for circulation on the shelf. These straits and gates (hereafter collectively referred to as
 288 gates when considering only the portion between the coast and 800 m isobath) are la-
 289 beled in Figure 1. The contour begins at the west side of Ellesmere Island (0 km), and
 290 the along-transect distance used in this paper is measured from that point counterclock-
 291 wise, first south along western Greenland then north along eastern Greenland. There-
 292 fore, along the contour, the gates are in order: (1) Davis Gate, (2) Cape Farewell Gate,
 293 (3) Sermilik Gate, (4) Denmark Gate, (5) Fram Gate, and (6) Nares Gate with the Nares
 294 Gate marking the eastern end of Nares strait and the end of the circum-Greenland con-
 295 tour. Between these gates we define the regional control volumes of the continental shelf
 296 as: (1) Northwest, (2) Southwest, (3) Narrow Shelf, (4) Wide Shelf, (5) Northeast, and
 297 (6) North, labeled in Figure 2. The Southeast region has been subdivided into the Nar-
 298 row and Wide sections because of differences in the cross-shelf exchange that we observed
 299 along the shelf break.

300 **3.2 Comparison to Observations**

301 The variables most critical for calculating volume and heat transports are veloc-
 302 ity and temperature. Since we are interested in eddy and mean processes, we consider
 303 both the strength of the mean currents and the spatial distribution and magnitude of
 304 the eddy kinetic energy (EKE). For temperature validation we focus on the temperatures
 305 in the Irminger Current as it is the major source of oceanic heat advected into our re-
 306 gion of interest. We look at the sea ice distribution on the southeast shelf in both mod-
 307 els. In addition, we compare the wind stress used to force each simulation.

313 The 2005-2009 climatology of speed averaged over the top 50 m of both models is
 314 shown in Figure 3. Both models show some evidence of the observed slope and coastal
 315 currents off southeast Greenland. The East Greenland Coastal Current has an offshore
 316 length scale of one Rossby radius but is present in both simulations, indicating that the
 317 resolution of the models is sufficient to capture some important mesoscale processes but
 318 not smaller-scale eddies. The currents strongly follow bathymetric contours, indicating
 319 the possibility that topographic steering plays a leading role in transport across the shelf
 320 break. North of Sermilik Trough (see Figure 1), between the coast and nearby troughs



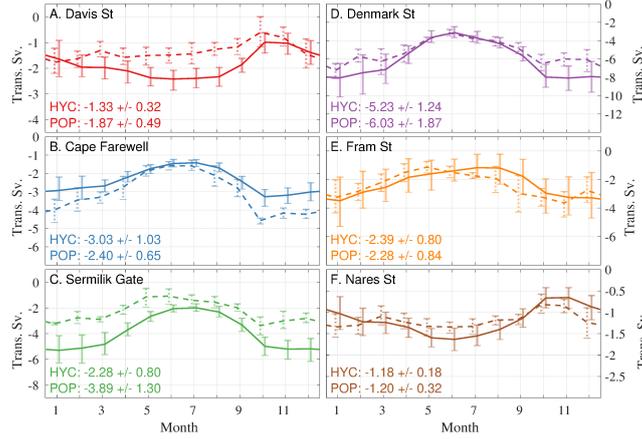
308 **Figure 3.** Average speed in the top 50 m over 2005-2009 for POP (A) and HYCOM (B).
 309 Average Eddy Kinetic Energy over the same period in the top 50 m over 2005-2009 for POP (C)
 310 and HYCOM (D). In (A,B) the red line shows the transect at 60°N. In (C,D) the boundary of
 311 two control volumes are shown in red: the interior Labrador Sea defined by the 2,000 m isobath
 312 and a box at the Denmark Strait.

321 there is a better defined coastal current in POP than in HYCOM. At 60°N, the Cape
 322 Farewell Gate marked by the red line in Figure 3, the peak velocity in HYCOM is 64 cm s⁻¹
 323 at a position 120 km from the coast. In POP there are two peaks in the surface speed:
 324 35 cm s⁻¹ located 97 km from the coast and 42 cm s⁻¹ located 155 km from the coast.
 325 See supplemental material for more details. This difference in current structure contributes
 326 to the difference in net transport onto the shelf between the two models (Figure 4). In
 327 general, the currents follow narrower pathways in POP than in HYCOM, possibly in-
 328 dicating less meandering. This is also supported by the higher EKE along the shelf break
 329 in HYCOM.

330 We calculate the EKE from the daily averages of velocity. We define

$$\text{EKE} = \frac{u'^2 + v'^2}{2} \quad (1)$$

331 with $u' = u - \bar{u}$, where u is the daily average velocity and \bar{u} is the monthly average of
 332 velocity. This defines eddies as anomalies that have a period between two days and one
 333 month. To obtain the near-surface expression of EKE, we use only the velocity in the
 334 top 50 m. The 2005-2009 climatology is plotted in Figure 3. In both models, west of Green-
 335 land there is an expanse of elevated EKE extending into the central Labrador Sea (out-
 336 lined in red in Figure 3 (C) and (D)). This represents an important pathway for the trans-
 337 port of heat and freshwater from the shelf into the Labrador Sea. The region of elevated
 338 EKE in POP starts at the shelf break, in contrast with HYCOM, which has elevated EKE
 339 on and off the shelf, possibly indicating a difference in the cross-shelf exchange between
 340 the two models in this region. EKE estimated from TOPEX/Poseidon satellite altime-
 341 try (Brandt et al., 2004) and surface drifters (Fratantoni, 2001) in this region shows a
 342 similar pattern of elevated EKE in the eastern Labrador Sea. The surface EKE from Brandt
 343 et al. (2004) in the West Greenland Current ranges from 400 to 800 cm² s⁻² for the pe-
 344 riod 1997-2001. Altimeter-based estimates are generally higher than those calculated here,
 345 which are based on velocity in the top 50 m. Estimates from 15 m drogued satellite-tracked
 346 surface drifter paths from 1990-1999 are 400 to 500 cm² s⁻², which is consistent with
 347 the maximum EKE of both models within the defined interior Labrador Sea control vol-
 348 ume (shown in red in Figure 3). In POP, the maximum EKE is 432 cm² s⁻² while in HY-
 349 COM it is 527 cm² s⁻². The average EKE in HYCOM is 53.2 cm² s⁻², with the 20th
 350 to 80th percentiles ranging from 20.0 to 62.9 cm² s⁻². POP values are similar: mean EKE
 351 is 51.1 cm² s⁻², with 20th to 80th percentiles from 10.1 to 82.0 cm² s⁻². The EKE of



365 **Figure 4.** Volume fluxes through transects defined in Figure 2 from POP (solid lines) and
 366 HYCOM (dashed lines) and error bars at the 20th and 80th percentiles; the annual mean and
 367 standard deviation are included on each plot. Transports are from (A) Davis Strait (B) Cape
 368 Farewell Gate (C) Sermilik Gate (D) Denmark Strait (E) Fram Strait and (F) Nares Strait; here
 369 strait refers to the entire transect between two land masses and gate refers to transport between
 370 the Greenland coast and the defined shelf break. Negative transport is southward.

352 both simulations has a similar magnitude and location to those observed in the Labrador
 353 Sea.

354 A second region of elevated EKE is a large region spanning from the northeast shelf
 355 at Scoresby Sound to where the Irminger Current retroflects south of the Denmark Strait.
 356 This corresponds to a region of large sea surface height anomalies observed by AVISO
 357 (Trodahl & Isachsen, 2018). Heightened EKE near the Denmark Strait is also consis-
 358 tent with observations of mesoscale eddies and boluses formed at the Denmark Strait
 359 overflow (Moritz et al., 2019). The average EKE in the defined box just south of the Den-
 360 mark Strait (outlined in red in Figure 3 (C) and (D)) is higher in POP ($133 \text{ cm}^2 \text{ s}^{-2}$)
 361 compared to HYCOM ($80.7 \text{ cm}^2 \text{ s}^{-2}$), and the maximum EKE in POP ($958 \text{ cm}^2 \text{ s}^{-2}$)
 362 is twice the maximum in HYCOM ($429 \text{ cm}^2 \text{ s}^{-2}$). In POP there is particularly strong
 363 band of EKE just south of the strait at the shelf break, while in HYCOM the maximum
 364 is broader and north of the strait.

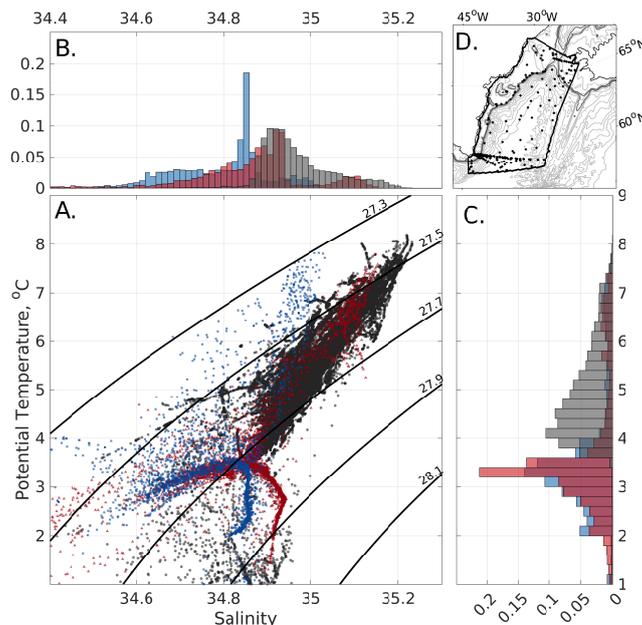
371 To further compare the velocity of the two models, we look at the net volume trans-
 372 port though the defined gates. For gates associated with major straits (Davis Strait, Den-
 373 mark Strait, Fram Strait, and Nares Strait), we compare the volume transport across

374 the entire strait to observational-based estimates. We find both the magnitude and sea-
375 sonal variability of volume transports to be consistent between the simulations. Figure
376 4 shows the mean seasonality for the transports in POP and HYCOM; negative values
377 indicate southward volume transport. Observational-based estimates of transports through
378 the same straits are shown in Table 1 with the 2005-2009 averages of both simulations.
379 Uncertainties of observational-based estimates of oceanic fluxes in Fram Strait have been
380 discussed (De Steur et al., 2018; Dushaw & Sagen, 2016) but apply broadly to straits
381 in this region that are poorly sampled in space and time. In addition, transport through
382 these straits varies interannually; the means presented may not be representative of the
383 true mean.

385 The transport through the Fram Strait is the primary source of Arctic Water in
386 the East Greenland Current. Both POP and HYCOM have estimates of the net volume
387 transport through the Fram Strait that are consistent with observations that include the
388 shelf currents. Volume transport through the Fram Strait has been shown to be sensi-
389 tive to both the inclusion of moorings that measure the shelf currents as well as the lat-
390 itude defining the strait (79°N vs 78.50°N). At the Denmark Strait, our model results
391 are compared to an estimate of the overflow transport (defined as $\sigma_\theta > 27.8 \text{ kg m}^3$), which
392 is lower than our estimates of the total transport through the strait, as expected. At the
393 Cape Farewell and Sermilik Gates, there are no constraining straits to use in defining
394 transports; therefore, we focus on the continental shelf volume fluxes from the coast to
395 the 800 m reference contour. At both gates there is similar magnitude and seasonality
396 to the volume fluxes. The winter maximum at Sermilik gate is weaker in HYCOM com-
397 pared to POP by roughly 1 Sv, but the summer minimum is similar in both models. Trans-
398 port at the Cape Farewell Gate has the opposite difference; the winter maximum in HY-
399 COM is 1 Sv greater than the maximum in POP. This is one indication that the cross-
400 shelf exchange in the Narrow Shelf region is different in the two simulations. The Cape
401 Farewell Gate is located at the same position as the Overturning in the Subpolar North
402 Atlantic Program (OSNAP) mooring array at 60°N; for both models, the maximum trans-
403 port is in winter, consistent with the observations of Le Bras, Straneo, Holte, and Hol-
404 liday (2018). At the Davis Strait, the average volume transport in HYCOM is about 0.54 Sv
405 less than in POP. Estimates of the volume transport from observations have a large range,
406 but the estimate from Curry et al. (2014) that covers a similar time period to our study

Table 1. Comparison of Observed and Modeled Volume Transports Through Straits

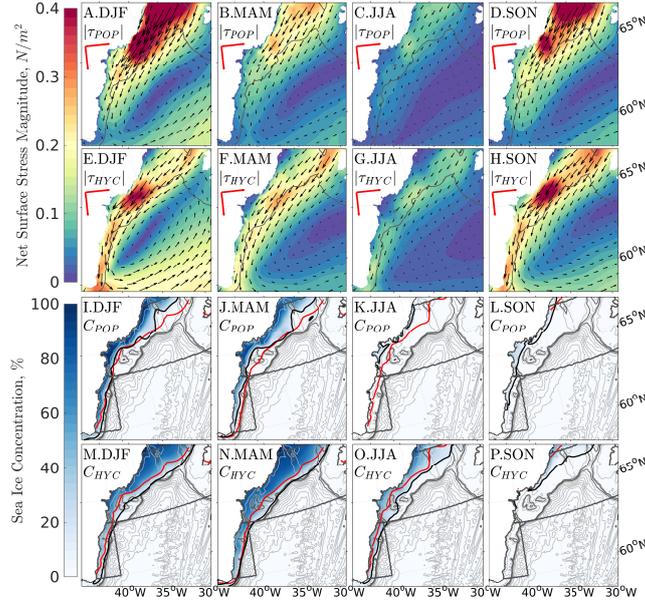
Strait	Time Period	Volume Transport	Reference
Fram Strait	1997-2000	-4 ± 2 Sv	(Schauer et al., 2004)
Fram Strait	1997-2007	-2.0 ± 2.7 Sv	(Beszczynska-Moeller et al., 2011)
Fram Strait	2005	-1.6 ± 3.9 Sv	(Tsubouchi et al., 2012)
Fram Strait	2005-2010	-2.39 ± 0.80 Sv	HYCOM
Fram Strait	2005-2010	-2.28 ± 0.84 Sv	POP
Denmark Strait Overflow	2007-2011	3.4 ± 1.4 Sv	(Jochumsen et al., 2012)
Denmark Strait	2005-2010	-5.23 ± 1.24 Sv	HYCOM
Denmark Strait	2005-2010	-6.03 ± 1.87 Sv	POP
Davis Strait	1987-1990	-2.6 ± 1.0 Sv	(Cuny et al., 2005)
Davis Strait	2004-2005	-2.3 ± 0.7 Sv	(Curry et al., 2011)
Davis Strait	2004-2010	-1.6 ± 0.5 Sv	(Curry et al., 2014)
Davis Strait	2005-2010	-1.33 ± 0.32 Sv	HYCOM
Davis Strait	2005-2010	-1.87 ± 0.49 Sv	POP
Nares Strait	Summer 2003	-0.8 ± 0.3 Sv	(Münchow et al., 2006)
Nares Strait	2003-2006	-0.57 ± 0.09 Sv	(Münchow & Melling, 2008)
Nares Strait	2005-2010	-1.18 ± 0.18 Sv	HYCOM
Nares Strait	2005-2010	-1.20 ± 0.32 Sv	POP



409 **Figure 5.** Comparison of WOA observations (gray), HYCOM (red), and POP (blue). (A) T-S
 410 diagram of observations from WOA (gray), HYCOM (red), and POP (blue). (B) Histogram of
 411 salinities from WOA, HYCOM, and POP, (C) same as (B) but for temperature. Locations of the
 412 WOA data are shown in (D) and limited to 200-800m depth.

407 has an average volume transport that falls within the range of both the POP and HY-
 408 COM simulations.

413 To validate the temperature and salinity properties of the models, we compared
 414 them to World Ocean Atlas (WOA) data collected in the eastern Irminger Sea. A con-
 415 trol volume (map in Figure 5(D)) was chosen to include part of the Irminger Sea and
 416 the southeast continental shelf. Observations collected between 2005-2009 were used (see
 417 supplemental material for details). The monthly mean that coincided with the timing
 418 of a given observation was used in order to limit the impact of eddy variability, and the
 419 model was sub-sampled at the closest grid point to the observation. We analyzed depths
 420 from 200 m to 800 m to limit the impact of surface variability. Both simulated and ob-
 421 served data sets were binned in temperature–salinity (Θ -S) space to calculate probabili-
 422 ty density functions (Figure 5). Overall, this comparison shows that the models are colder
 423 and fresher than the observations. The mean temperature is $3.6 \pm 1.4^\circ\text{C}$ in POP, $3.5 \pm 1.4^\circ\text{C}$
 424 in HYCOM, and $4.6 \pm 1.2^\circ\text{C}$ in WOA, where ranges given here are one standard devi-



427 **Figure 6.** Seasonal averages of daily net surface stress fields from CORE-II CIAF from 2005-
 428 2009 used to force the 0.1° simulation (A-D) and CFSR forcing from 2005-2009 used for the
 429 HYCOM 0.08° simulation (E-H). Seasonal averages of sea ice concentration from 0.1° POP (I-
 430 L) and 0.08° HYCOM (M-P). The ice edge defined by the 15% concentrations from the model
 431 simulations are shown in black, and from SSM/I in red (25 km resolution).

425 ation. Similarly, POP is the freshest of the three data sets with a mean salinity of 34.8 ± 0.13 PSU,
 426 compared to 34.8 ± 0.20 PSU in HYCOM, and 35.0 ± 0.08 PSU in WOA.

432 Over the East Greenland Continental Shelf, persistent northerly winds are down-
 433 welling favorable and play a significant role in setting up the currents on the shelf (Sutherland
 434 & Pickart, 2008). Where present, sea ice mediates how wind stress drives the ocean. Sea
 435 ice reflects more solar radiation than the ocean and limits the radiative transfer into the
 436 ocean, limiting the surface heat flux reaching the ocean surface. The northerly winds are
 437 strongest in the winter when sea ice extent is largest; the seasonal magnitude of the net
 438 surface stress driving each model is shown in Figure 6.

439 The sea-ice edge position is determined in part by the meandering and strength
 440 of the currents on the shelf. In Figure 3, the topographic steering at Kangerdlugssuaq
 441 Trough is more pronounced in POP. Similarly the sea-ice edge somewhat follows the to-
 442 pography of the trough. This difference in the two simulations is also pronounced at the
 443 Sermilik Troughs region, where the sea ice is usually shoreward of the troughs in POP

444 but in HYCOM covers the entire shelf. We compare the total ice area in two boxes (black
 445 boxes in Figure 6), one covering the Narrow Shelf region and the other the Wide Shelf
 446 region. The total area covered by sea-ice in winter (December, January and February)
 447 in the SSM/I data is 0.6×10^6 km² in the Narrow Shelf box and 5.5×10^6 km² in the Wide
 448 Shelf box. In HYCOM and POP the winter sea-ice area in the Wide Shelf box is 5.9×10^6 km²
 449 and 3.8×10^6 km²; in the Narrow Shelf box it is 1.0×10^6 km² and 0.7×10^6 km², respec-
 450 tively. Sea-ice area over the Wide Shelf box is consistently overestimated in HYCOM
 451 and underestimated in POP, with the exception of fall (September, October, and Novem-
 452 ber) when POP has 60% more sea-ice area than SSM/I. In the Narrow Shelf box there
 453 is a less clear difference in the simulations and observations, both simulations have more
 454 fall and winter sea-ice by area than SSM/I. In the spring (March, April, and May) both
 455 models underestimate sea-ice area (47% in POP, 20% in HYCOM). For POP this con-
 456 tinues into summer (June, July and August, 88% less than SSM/I), while in HYCOM
 457 it reverses (43% more than SSM/I).

458 The net surface stress includes both the wind stress and the stress from the sea ice
 459 onto the ocean. Over the shelf, the location and spatial structure of the strongest sur-
 460 face stress differ between the forcings. In the POP surface forcing, the maximum stress
 461 is over the Wide Shelf with two local maxima: one to the south of Kangerdlugssuaq Trough
 462 and the other to the north of Denmark Strait. In winter, these maxima are strongest and
 463 merge into a single region of high surface stress over the majority of the Wide Shelf re-
 464 gion. The POP surface stress in the Narrow Shelf region is weaker than the stress over
 465 the Wide Shelf. The HYCOM surface forcing has a maximum south of Kangerdlugssuaq
 466 Trough that is similar in magnitude to the maximum in POP. However, the maximum
 467 to the north of the Denmark Strait is weaker and further offshore than the northern max-
 468 imum in POP. The HYCOM surface stress along the Narrow Shelf and near Cape Farewell
 469 is higher and more coastally trapped than the surface stress in POP. There are mod-
 470 erate winds all along the southeast shelf slope throughout the year that are strongest at
 471 Cape Farewell.

472 The differences in sea-ice cover and sea-ice–ocean stress are not sufficient to explain
 473 the difference in the net surface stress between the two simulations in this region. We
 474 attribute the majority of the difference to discrepancies in the CORE2-IAF and CFSR
 475 wind products. These differences in wind stress would also drive sea-ice dynamics dif-
 476 ferently, thus compounding the difference in net surface stress. Overall, the surface stress

477 is stronger in POP over the Wide Shelf region; in HYCOM the region of high net sur-
 478 face stress is weaker but extends farther south to Cape Farewell. These differences could
 479 be a primary driver of differences in cross-shelf transports and the shelf currents.

480 Overall, our models compare well with each other and with observations. The res-
 481 olution allows for shelf break currents with realistic strength and seasonality. The trans-
 482 ports through major straits are consistent among the two models and observations. The
 483 water masses present in both models and observations are similar with moderate tem-
 484 perature biases in both models. The wind stresses used to force each simulation have sig-
 485 nificant differences in spatial structure over the southeast Greenland Shelf but similar
 486 seasonal cycles.

487 3.3 Heat Transport Definition

488 Volume fluxes are calculated using daily means of velocity. The net volume trans-
 489 port across the 800 m isobath is defined as

$$V_{800} = \int_H^0 \int_0^L \hat{v} d\hat{x} dz, \quad (2)$$

490 where \hat{x} is the along-boundary direction and \hat{v} is the velocity component perpendicu-
 491 lar to the transect. In the case of the 800 m isobath, the normal direction, n , is defined
 492 such that $\hat{v} = v\hat{n}$ is onto the shelf. In the case of the gates, the volume flux is calcu-
 493 lated similarly, but the normal direction is northward. This allows us both to look at the
 494 overall volume flux onto the shelf and to construct budgets for the individual shelf re-
 495 gions by considering whether the gate is at the northern or southern boundary of the re-
 496 gion. (If the gate is the northern boundary, the normal direction must be reversed to point
 497 into the box.) Both models are volume conserving, with small free surface variations. There-
 498 fore we can use volume budgets to validate our transport estimates.

499 Heat flux is calculated using daily means of potential temperature and velocity from
 500 both models. For the heat transport across the 800 m isobath we define

$$\Phi_{800} = \int_H^0 \int_0^L \rho c_p (\theta - \theta_{ref}) \hat{v} d\hat{x} dz, \quad (3)$$

501 where ρ is the density of seawater, c_p is the specific heat capacity of seawater, and θ is
 502 the reference potential temperature. This definition is used both for the transport across
 503 the 800 m isobath (Φ_{800}), and through the various gates (Φ_G). As with the volume trans-

504 port, positive volume flux is onto the shelf, and gate fluxes are positive in the northern
 505 and eastern directions.

506 Changes in the heat content of a specified volume are given by the heat fluxes in-
 507 tegrated over the bounding surface. Here, the change in heat of a given volume is de-
 508 fined by the convergence of the lateral ocean fluxes, the net surface heat flux, and heat
 509 storage (change in heat content of the control volume):

$$\rho c_p \frac{d\Theta_{vol}}{dt} = (\Phi_{800} + \Phi_{G-north} + \Phi_{G-south}) + \Phi_{SHF}. \quad (4)$$

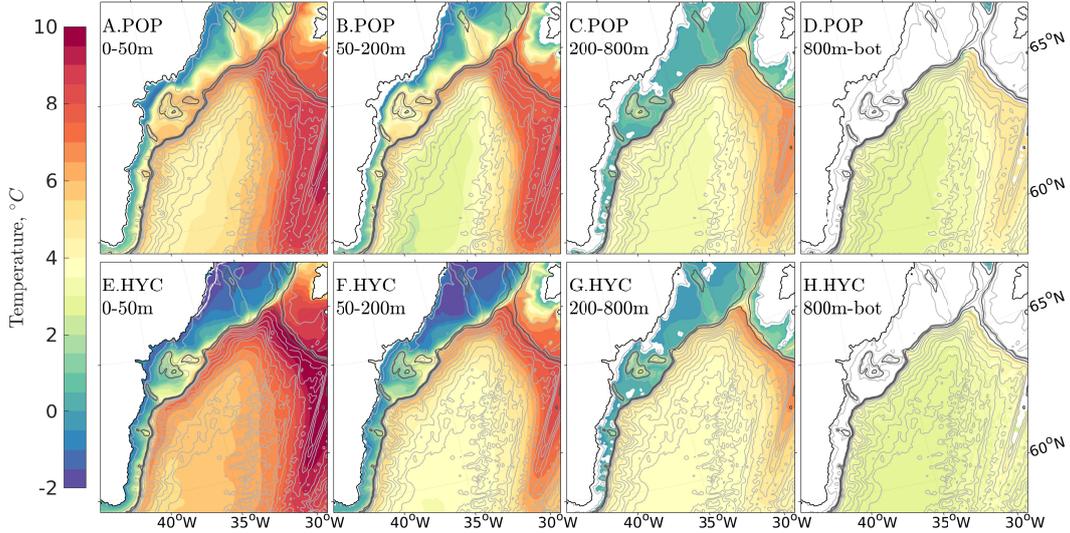
510 Without having saved the heat storage term as the model ran, we cannot close the heat
 511 budget exactly. Using daily averages we can approximate the change in heat content over
 512 time, but acknowledge that there is some variability that is obscured by the daily av-
 513 eraging. Surface heat fluxes from the HYCOM simulations were not saved. The choice
 514 of reference temperature does not change the net heat transport into an enclosed region
 515 (Bacon & Fofonoff, 1996; Schauer & Beszczynska-Möller, 2009). We have used a refer-
 516 ence temperature $\theta_{ref} = -1.8^\circ\text{C}$, which is the salinity-independent freezing tempera-
 517 ture in POP (R. Smith & Gent, 2002).

518 4 Results

519 4.1 Heat Content of the Southeast Shelf

524 Many tidewater glaciers terminate in the Southeast region, making the region im-
 525 portant to ocean-driven melting of the GIS (Millan et al., 2018). We specifically exam-
 526 ine the vertical structure of the average temperature, shown in Figure 7 in four layers:
 527 surface (0-50 m), mid-depth (50-200 m), shelf bottom (200-800 m), and deep (800 m to
 528 bottom). On the shelf, the importance of bathymetry in controlling warming is appar-
 529 ent as warm water is present in areas with deeper bathymetry. The two main regions
 530 of warm water on the shelf at all depths are in the Kangerdlugssuaq and Sermilik Troughs.

531 The HYCOM shelf water is colder than that in POP, with average temperatures
 532 in the Wide Shelf region of $0.95 \pm 0.55^\circ\text{C}$ versus $2.15 \pm 0.64^\circ\text{C}$, respectively (see Table 2
 533 for average temperatures in each control volume). The position of the warm-cold front
 534 in the upper layers is primarily along the 800 m isobath in HYCOM (E and F) but is
 535 further onshore in POP (A and B). Warm water incursions onto the shelf over deep bathymetry,
 536 such as Kangerdlugssuaq Trough and Sermilik Deep, in the surface layer are stronger in

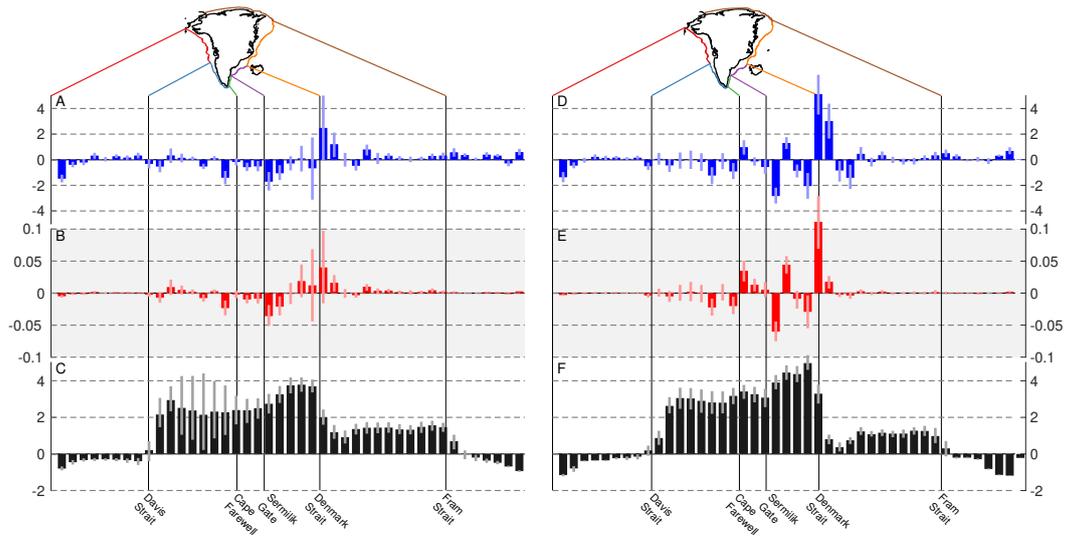


520 **Figure 7.** Climatology from 2005-2009 of potential temperature in POP (A-D) and HYCOM
 521 (E-H) for depth ranges: 0-50 m (A,E), 50-200 m (B,F), 200-800 m (C,G), and 800 m–bottom
 522 (D,H). Bathymetry contours are shown in light gray, with the 500 m and 800 m isobaths in dark
 523 gray.

537 POP (A and B) compared to HYCOM (E and F). This is consistent with the topograph-
 538 ically steered current speed observed in the top 50 m (Figure 3); note the difference in
 539 bathymetry at the mouth of the trough (supplemental material). This difference is re-
 540 duced in the 200-800 m layer, (C for POP, G for HYCOM) as are differences in the Irminger
 541 Basin temperature. The water within the Irminger Basin is generally warmer in HYCOM
 542 than in POP; here the basin is defined as the region deeper than the 800 m isobath and
 543 north of 60°N, the region enclosed by the black contour in Figure 7. The average basin
 544 temperature in the surface layer of HYCOM (E) is 6.9°C; the average in POP (A) is 6.1°C.
 545 This difference between the two models is not present from 50 - 800 m, where the basin
 546 temperatures differ by $<0.1^{\circ}\text{C}$, but from 800 m to the bottom, POP (D) is 0.5°C warmer
 547 than HYCOM (H).

548 4.2 Average Transports onto the Greenland Shelf

556 Figure 8 shows the 5-year climatology of the vertically integrated volume (A for
 557 POP and D for HYCOM) and heat (B for POP and E for HYCOM) fluxes in each sim-
 558 ulation in 100 km sections of the contour. The gates are marked along the plot to show
 559 which regions have the strongest fluxes and warmest temperature. The average temper-



549 **Figure 8.** Bar graph of net volume fluxes (Sv, blue), net heat fluxes (PW, red), and average
 550 temperature ($^{\circ}\text{C}$, black) in 100 km sections of the transect for POP (A-C) and HYCOM (D-F).
 551 For both volume and heat fluxes, positive values indicate flux onto the shelf. Dark bars are the 5
 552 year average from 2005-2009, with light bars representing the 20th and 80th percentile range. In
 553 both models, the strongest on shelf fluxes are near the Denmark Strait. In POP this maximum is
 554 associated with strong variability; in HYCOM the heat flux is consistently onto the shelf at this
 555 location.

Table 2. Average Transports Through Gates and Across the Shelf.

Section	Length (km)		V (Sv)		Φ (TW)		T_{avg} ($^{\circ}$ C)	
	POP	HYCOM	POP	HYCOM	POP	HYCOM	POP	HYCOM
Northwest Shelf	1,785	1,864	-1.34 \pm 0.51	-1.46 \pm 0.44	-6.11 \pm 5.86	-5.54 \pm 4.76	-0.27 \pm 0.18	-0.65 \pm 0.15
Davis Gate	166	192	0.25 \pm 0.74	-0.27 \pm 0.59	8.52 \pm 13.2	-7.85 \pm 9.73	-	-
Southwest Shelf	1,622	1,651	-2.40 \pm 0.79	-2.15 \pm 0.84	-29.9 \pm 10.4	-16.6 \pm 16.0	1.96 \pm 1.36	1.96 \pm 0.84
Cape Farewell Gate	77	62	-2.62 \pm 0.75	-2.45 \pm 0.85	-38.7 \pm 19.8	-33.4 \pm 18.6	-	-
Narrow Shelf	503	503	-1.13 \pm 0.62	-0.59 \pm 0.45	-18.0 \pm 13.9	28.2 \pm 15.7	1.65 \pm 1.02	1.41 \pm 0.81
Sermilik Gate	74	87	-3.53 \pm 1.20	-3.03 \pm 1.12	-53.2 \pm 21.6	-39.5 \pm 21.9	-	-
Wide Shelf	1,021	994	-0.90 \pm 0.66	0.77 \pm 0.76	16.4 \pm 13.8	55.0 \pm 23.3	2.15 \pm 0.64	0.95 \pm 0.55
Denmark Gate	257	361	-4.07 \pm 1.42	-2.28 \pm 0.95	-33.2 \pm 18.9	-13.0 \pm 8.43	-	-
Northeast Shelf	2,315	2,303	3.07 \pm 1.11	1.63 \pm 0.90	44.2 \pm 19.6	16.1 \pm 8.17	0.11 \pm 0.37	-0.62 \pm 0.17
Fram Gate	298	266	-0.99 \pm 0.48	-0.62 \pm 0.23	-1.84 \pm 1.57	-0.94 \pm 0.60	-	-
North Shelf	1,314	1,317	1.96 \pm 0.31	1.51 \pm 0.33	5.65 \pm 1.73	1.13 \pm 1.61	0.49 \pm 0.40	-0.99 \pm 0.03
Nares Gate	76	70	-1.15 \pm 0.40	1.19 \pm 0.27	-3.60 \pm 1.45	3.77 \pm 0.83	-	-

560 ature at the shelf break in each section is shown in Figure 8 C (POP) and F (HYCOM).
561 On the whole the two simulations have a similar pattern of volume and heat fluxes. In
562 both, the strongest on shelf flux is near the Denmark Strait, with weak on shelf flux north
563 of the strait and mostly off shelf flux in western Greenland. The magnitudes of the fluxes
564 and their variability differ between the two simulations. HYCOM has warmer water (panel
565 F), with less variability in temperature compared to POP (panel C). Combined with stronger
566 volume fluxes in HYCOM (panel D vs panel A) the result is greater magnitude heat fluxes
567 in HYCOM (panel E vs B). While the HYCOM simulation does not have the same tem-
568 poral variability as POP, there is along-transect variability where regions of strong off-
569 shelf flux are adjacent to those with strong on-shelf flux. In POP, the temporal variabil-
570 ity (shown here by the relative size of the 20th to 80th percentile range) is large rela-
571 tive to the mean between Sermilik Gate and the Denmark Strait. Topographic Rossby
572 waves in this region could explain some of this variability, as discussed further in Sec-
573 tion 4.3.

574 Table 2 summarizes the net heat and volume fluxes through every section and gate.
575 The results from both HYCOM and POP simulations show that temperature within a
576 given control volume is not closely linked to cross-isobath heat transport. This implies
577 that shelf circulation and surface heat fluxes are important to the regional heat budget
578 of the Greenland Continental Shelf. The results in Table 2 also indicate the role of the
579 shelf circulation through the gate fluxes. The heat flux through the gates along the east
580 coast of Greenland is southward. In both POP and HYCOM, there is less heat flux at
581 Cape Farewell than at the Sermilik Gate indicating that the Narrow Shelf is a region of
582 heat loss, despite it being a region of net on-shelf heat flux in HYCOM. Note Figure 4
583 shows the volume transport across the entire straits and should not be compared to these
584 transports, which only include transport from the coast to the shelf contour.

585 In west Greenland, we expect to see off-shelf volume and heat transport in agree-
586 ment with previous studies (e.g., Dukhovskoy et al., 2019; Böning et al., 2016; Schulze Chre-
587 tien & Frajka-Williams, 2018; Myers et al., 2009). Both simulations are consistent in this
588 region, with weak seasonal cycles of heat and volume transport. At the Davis Strait, the
589 northward heat flux is small, and the northwest control volume is much colder than the
590 southwest volume. The volume-averaged shelf temperature of the Southwest region is
591 highly variable, and the fall peak is the warmest volume-average temperature of any re-
592 gion. The presence of warm ocean water in this region is consistent with observations

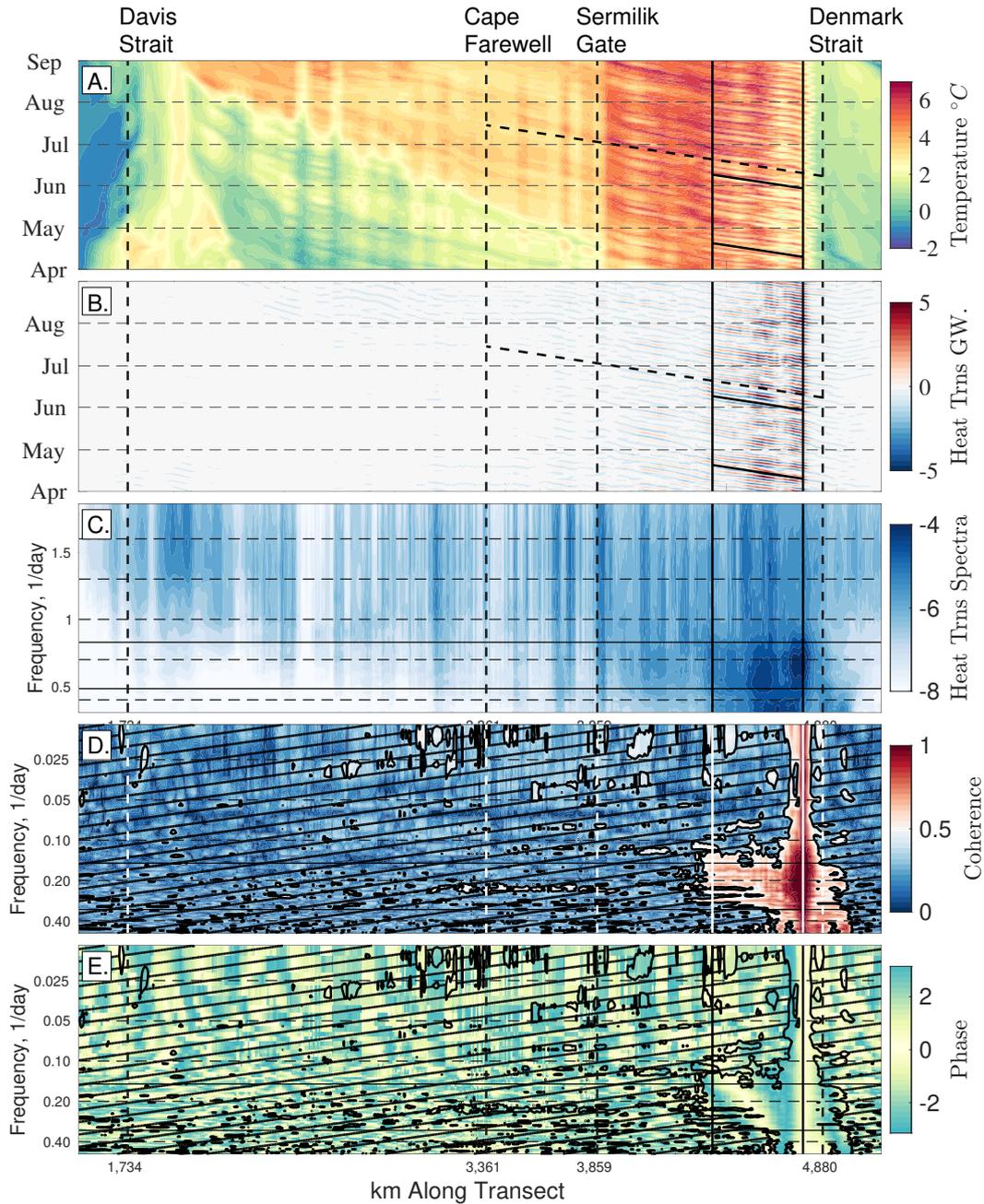
593 of ocean-driven melting of the ice sheet in west Greenland (Holland et al., 2008). Cross-
 594 isobath heat transport is negative in the Southwest region, consistent with the source
 595 of heat to this region originating from southward heat transport at Cape Farewell or sur-
 596 face heat fluxes. Correlation between the heat flux at the Cape Farewell Gate and heat
 597 content in the Southwest region is 0.87 in POP and 0.74 in HYCOM; both are signif-
 598 icant at a 95% confidence level. Using the surface heat flux time series saved from the
 599 POP simulation, we find that the net surface heat flux and heat content in the South-
 600 west region are out of phase, resulting in low correlation. In both models, heat trans-
 601 port through the Cape Farewell Gate as well as the shelf temperature peak in the fall;
 602 in POP the net surface heat flux is maximum in the summer.

603 **4.3 Daily Variability of Transport along the Greenland Shelf**

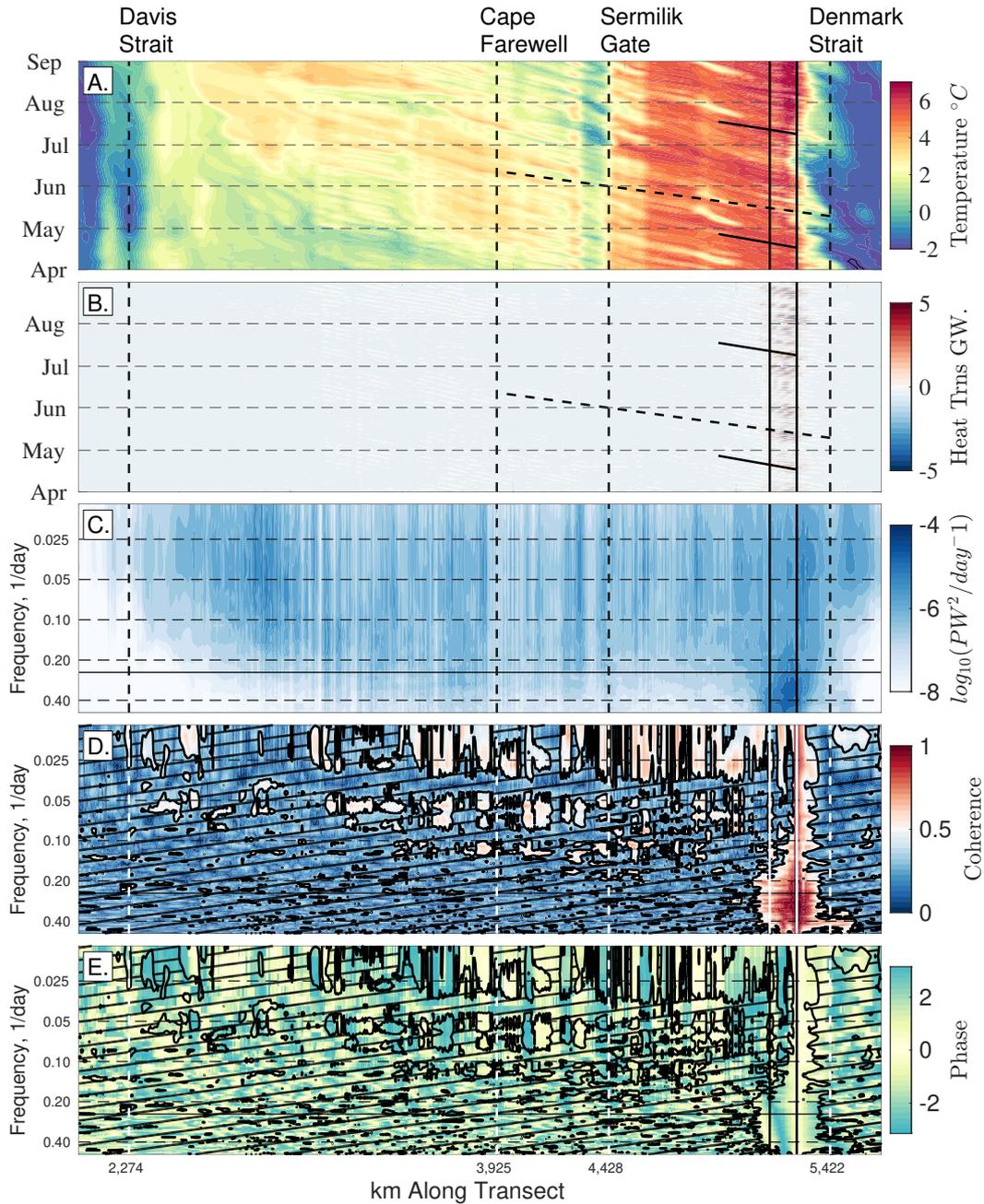
622 Heat transport is also highly variable along the transect within the defined shelf
 623 regions and on multi-day timescales. Figures 9A and 10A show 5-month-long Hovmöller
 624 diagrams of temperature at 200 m in 2005 for each model at the shelf break, illustrat-
 625 ing the seasonal progression of warm water from Denmark Strait to Davis Strait. To re-
 626 duce noise in all variables from currents meandering across the isobath, a 50 km box-
 627 car filter is applied. Hovmöller diagrams of 200 m temperature for the full five-year pe-
 628 riod are included in the supplemental material.

629 At the Denmark Strait there is a front between the cold Norwegian Seas water and
 630 the warm Atlantic water in the Irminger Current in both models, but the north part of
 631 the front is much colder in HYCOM (Figure 10A) consistent with the average shelf tem-
 632 peratures in both simulations ($-0.62 \pm 0.17^\circ\text{C}$ compared to $0.11 \pm 0.37^\circ\text{C}$ in POP, see Ta-
 633 ble 2). The warmest water at the shelf break in both models is along the Wide Shelf re-
 634 gion (between the Denmark and Sermilik Gates) and does not vary much seasonally. In
 635 POP between Sermilik Gate and Cape Farewell seasonal warming occurs in May (Fig-
 636 ure 9). However, in HYCOM (Figure 10) the temperature over this portion of the shelf
 637 break shows more high-frequency variability than seasonal change. These differences are
 638 consistent with the annual cycles of temperature in the Southwest region and heat trans-
 639 port through the Cape Farewell and Davis Gates.

640 In both models, there is a high-frequency signal generated at or intersecting the
 641 shelf break south of the Denmark Strait in roughly the same location as the cold-warm



604 **Figure 9.** POP results showing:(A) Hovmöller diagrams from April to September 2005 of
 605 temperature at 200 m, (B) vertically integrated heat transport with a 3-7 day band pass filter,
 606 (C) spectra of heat transport at each location along the contour with horizontal lines showing
 607 the frequency band that was used to produce (B), (D) coherence between heat transport at every
 608 location and 90 km south of the Denmark Strait, and (E) the associated phase. Vertical dashed
 609 lines show the locations of the gates, and solid vertical lines show the region of the propagating
 610 signal from 90 to 500 km south of the Denmark Strait. Cross-hatching of coherence and phase
 611 indicates that the coherence is not significant. Error for the spectra are estimated using a χ^2
 612 distribution with a 95% significance level such that the range between high and low error estimates
 613 is $\log_{10}(0.6)$.



614 **Figure 10.** As in Figure 9. HYCOM results showing:(A) Hovmöller diagrams from April to
 615 September 2005 of temperature at 200 m, (B) vertically integrated heat transport with a 2-6
 616 day band pass filter, (C) spectra of heat transport at each location along the contour with hor-
 617 izontal lines showing the frequency band that was used to produce (B), (D) coherence between
 618 heat transport at every location and 150 km south of the Denmark Strait, and (E) the associated
 619 phase. Vertical dashed lines show the locations of the gates, and solid vertical lines show the
 620 region of the propagating signal from 150 to 270 km south of the Denmark Strait. Cross-hatching
 621 of coherence and phase indicates that the coherence is not significant at the 95% level.

642 front (Figures 9A and 10A). In POP (Figure 9), the origin of these signals is consistently
 643 90 km south of the Denmark Strait. In HYCOM (Figure 10) the position of the cold-
 644 warm front meanders and changes in strength over the months shown. These high-frequency
 645 signals are generated regularly throughout the year, including months not shown here.
 646 As these signals propagate along the transect they result in extreme high and low heat
 647 transports.

648 Figures 9B and 10B show the band-pass filtered vertically integrated heat trans-
 649 port, and Figures 9C and 10C show the spectra of the vertically integrated heat trans-
 650 port. In both models the heat transport spectra have peaks at high frequency south of
 651 the Denmark Strait. In POP (Figure 9), there are three localized regions of high-frequency
 652 variability, two with a frequency of about 0.30 day^{-1} , and one of higher frequency. In
 653 HYCOM (Figure 10) the high-frequency peak is more localized from 0.24 day^{-1} to 0.5 day^{-1} ,
 654 the Nyquist frequency. Therefore, to isolate the heat transport associated with these prop-
 655 agating signals, the models were band-pass filtered with different ranges: for POP the
 656 range is a period of 3-7 days, and for HYCOM the range is 2-4 days. The band-pass fil-
 657 tered heat transport in both models (Figures 9(B) and 10(B)) shows a propagating sig-
 658 nal, though the signal travels only 120 km in HYCOM, while in POP it continues for
 659 410 km. The location where the signal dissipates in HYCOM (Figure 10) coincides with
 660 the mouth of Kangerdlugssuaq Trough. In POP (Figure 9) the signal dissipates on the
 661 north end of the Sermilik Troughs. In both cases the dissipation or on/off-shelf shifting
 662 of the signal occurs where there is a change in bathymetry. In both models (Figures 9C
 663 and 10C), the high-frequency energy in the spectra of vertically integrated heat trans-
 664 port decays southward along the shelf. From Sermilik Gate to Davis Strait, much of the
 665 variability in both models is at frequencies below 0.15 day^{-1} . In HYCOM (Figure 10),
 666 the energy in that range is consistent over that portion of the shelf break; however, in
 667 POP (Figure 9) there are regions of very low spectral energy especially where the winds
 668 are weakest.

669 The coherence of heat transport time series at each location along the contour and
 670 the heat transport at the origin location of the signals, with the associated phase, was
 671 used to assess these signals; Figures 9D and 10D show the coherence squared, and Fig-
 672 ures 9E and 10E show the associated phase. Where the coherence squared is less than
 673 the 95% confidence level $\Gamma_{XY}^2 = 0.39$, coherence and phase are hatched. For both mod-
 674 els, there are regions of strong coherence both north (upstream, closer to the Denmark

675 Strait) and south (downstream, farther from the Denmark Strait). The upstream coher-
676 ence shows the possible origin of the signal. In HYCOM (Figure 10) the coherence is sig-
677 nificant north of the Denmark Strait in the same narrow high-frequency band (0.24-0.5 day⁻¹).
678 In POP (Figure 9) the coherence is significant north of the Denmark Strait across most
679 frequencies in the 3-7 day band. The fact that the forcing differs could potentially ex-
680 plain the differences in the coherent responses. Both the model bathymetry and the choice
681 of shelf break contours could also contribute to the difference. In both models, where the
682 coherence is significant south of the Denmark Strait, the phase shows evidence of a prop-
683 agating signal. In HYCOM, there is also coherence at low frequencies along the shelf to
684 Cape Farewell and partway along the western Greenland shelf. This coherent signal is
685 not present in POP and could be related to the larger scale meandering of the front where
686 the signal originates, or any of the differences in model set up listed above that could
687 explain differences in the high frequency signals, but thorough analysis is beyond the scope
688 of this study.

689 As in Figures 9C and 10C, the frequencies where the signals are coherent along the
690 transect (from right to left in Figures 9 and 10) differ between simulations. From the phase,
691 we can estimate the phase velocity of the signal (Münchow et al., 2020; Pickart & Watts,
692 1990). A middle frequency of each band of coherence was used, $f_{POP}=0.2$ day⁻¹ for POP
693 and $f_{HYCOM}=0.34$ day⁻¹ for HYCOM. A location was chosen along the contour near
694 where the coherence at that frequency is no longer significant, 500 km south of the Den-
695 mark Strait in POP, 270 km in HYCOM; the distance between the two locations is D .
696 At that frequency and location, the phase is $\Theta_{XY} = 128^\circ$ in POP and $\Theta_{XY} = 52^\circ$ in
697 HYCOM. We calculate the phase speed as $c_p = f(360/\Theta_{XY})(D/\cos \Delta)$, where Δ is
698 the angle between the wavenumber vector and the direction of the shelf break. In this
699 region, the difference in the direction of the average velocity and shelf break direction
700 was used to estimate Δ , with $\Delta = 4^\circ$ in POP and $\Delta = 14^\circ$ in HYCOM. For POP,
701 the resulting phase velocity is $c_p = 2.6$ m/s and wave length $\lambda = 1,155$ km; for HY-
702 COM, the resulting phase velocity is $c_p = 3.3$ m/s and wave length $\lambda = 863$ km. The
703 spectra, coherence, and phase used for estimating the phase velocity are shown in the
704 Supplemental Material.

705 In addition, we use the dispersion relation for topographic Rossby waves

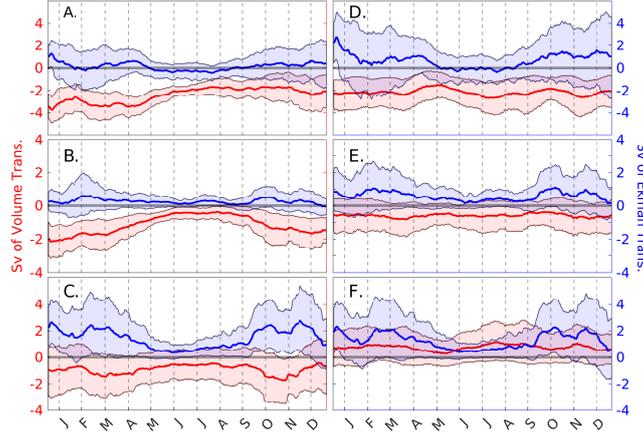
$$\omega = \frac{N\alpha \sin(\phi)}{\tanh(KL_D)} \quad (5)$$

706 to estimate the orientation of the wavenumber vector, $\phi = k/K$ where $\vec{k} = (k, l)$ and
 707 $K = \sqrt{k^2 + l^2}$. The stability frequency N is estimated from profiles of the 5-year av-
 708 erage stratification in this region of the shelf, $N = 0.003 \text{ s}^{-1}$ in POP and 0.002 s^{-1}
 709 in HYCOM. The average bottom slope estimated from the change in local bathymetry
 710 to be $\alpha = 0.02$ for POP and HYCOM. The Rossby radius of deformation $L_S = ND/f$
 711 is calculated at each point along the contour, using the local depth and Coriolis param-
 712 eter, the average value of $L_D = 12 \text{ km}$ in this region of both models. For POP, using
 713 this equation and the estimated wavelength from the phase analysis, the value of $\phi =$
 714 1° , compared to the orientation of the current ellipse at the shelf break $\theta = -6^\circ$. In
 715 HYCOM, the angle of the wavenumber vector is $\phi = 4^\circ$, and the average direction ori-
 716 entation of the current ellipse is $\theta = -0.5^\circ$. This agreement between the direction of
 717 wave propagation estimated from the topographic Rossby waves dispersion relationship
 718 and the velocity ellipses at the shelf break is consistent with the hypothesis that these
 719 waves are topographic Rossby waves.

720 These waves are present in both the vertically integrated heat transport and the
 721 200 m temperature. In the previous section, we showed that offshore Ekman transport
 722 was not the leading mechanism of heat transport in this region. Because of the strength
 723 of the propagating signal across all depths in this region, we hypothesize it to be the pri-
 724 mary mechanism of cross-shelf exchange. The mean heat transport in this region in Fig-
 725 ure 8 over the five years of our study period is highly variable but has a net on-shelf trans-
 726 port. Therefore, we conclude that these waves not only perturb the thermocline but also
 727 result in a net heat transport onto the shelf.

728 5 Discussion

729 Our goal is to better understand the heat budget of the shelf, and the mechanisms
 730 governing heat transport. From the observed heat transports, we find variability on two
 731 timescales. On the seasonal scale, warming of the shelf is associated both with cross-shelf
 732 transports and with the transport of heat by currents on the shelf. On shorter time scales,
 733 there is evidence that topographic Rossby waves propagate from south of the Denmark
 734 Strait to locations along the shelf where the bathymetry changes rapidly. For the sea-
 735 sonal variability of volume and heat transport, the role of winds through Ekman trans-
 736 port is proposed as a leading mechanism. For the waves, the local and regional winds
 737 are investigated as a source of the high-frequency signals. Together, these mechanism



743 **Figure 11.** Comparison of the mean of the daily Ekman transport (Sv, blue) to net cross-
 744 isobath volume transports (Sv, red) from 2005-2009 in three focus regions: Southwest (A,D),
 745 Narrow Shelf (B,E), and Wide Shelf (C, F). (A-C) are results from POP and CORE-IAF, (D-F)
 746 are results from HYCOM and CFSR. Time series are smoothed with a 30-day running mean;
 747 shaded areas show the 20th and 80th percentile range for each day. Transports are positive shore-
 748 ward.

738 set the boundary conditions of the continental shelf heat budget. The combination of
 739 both mechanisms controls the heat flux in the shelf region where there is strongest shore-
 740 ward heat flux. The connection between the mechanisms and the heat budget is essen-
 741 tial to understanding shelf temperature sensitivity to large scale changes.

742 5.1 Mechanism One: Ekman Dynamics

749 Winds in southeast Greenland are predominantly downwelling favorable near the
 750 coast and along the shelf break. Northerly winds drive Ekman transport that advects
 751 surface water onto the shelf setting up a strong sea surface gradient and driving down-
 752 welling, forcing deeper water off the shelf. Using the net surface stress (including both
 753 winds and sea ice), we can estimate the resulting Ekman transport across our selected
 754 transect. We define the net Ekman downwelling transport ($\text{m}^3 \text{s}^{-1}$) across a given sec-
 755 tion of the 800 m isobath as

$$V_{DW} = \int_0^L \frac{1}{f\rho} \tau_{along} dx, \quad (6)$$

756 where ρ is the density of sea water (not property dependent), and f is the Coriolis pa-
 757 rameter, which varies with latitude.

758 We also consider whether differences in the two forcing products (CORE-IAF for
 759 POP and CFSR for HYCOM) may influence the transport in both simulations. We ex-
 760 pect that, in two atmospherically forced simulations with robust physics, transport mech-
 761 anisms would agree; however, we must consider how differences in the forcing could pro-
 762 duce different dynamics. To quantify this, in the most basic way, we use linear regres-
 763 sion to compare the difference of the Ekman transports $V_{DW\ diff} = V_{DW-HYCOM} -$
 764 V_{DW-POP} and the difference in volume transports $V_{800\ diff} = V_{800-HYCOM} - V_{800-POP}$.

765 In Figure 11, time series of V_{DW} and V_{800} are plotted for three different regions:
 766 Southwest, Narrow Shelf, and Wide Shelf; positive transports are onto the shelf. The dif-
 767 ferences in V_{DW} estimated from POP and HYCOM highlight the differences in the wind
 768 forcing. The mean annual cycle illustrates the seasonality of the wind forcing and cross-
 769 isobath transports, with the acknowledgement that by using only five years of data, year-
 770 to-year variability is still present. Vertical profiles of volume fluxes in the Narrow and
 771 Wide Shelf regions have shown that while the net volume transport is off the shelf, there
 772 is onshore flux near the surface (upper 250 m). In the Southwest region, the volume trans-
 773 port is generally off the shelf over the entire water column.

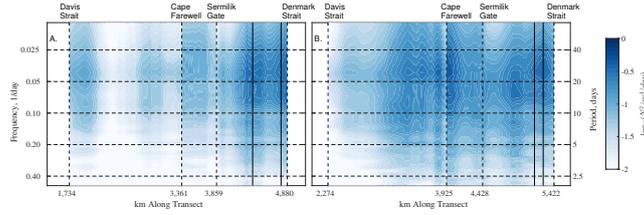
774 In the Southwest region (panels A and D), the surface forcing in both models drives
 775 seasonally varying on-shelf (winter) and off-shelf (summer) Ekman transports. Dukhovskoy
 776 et al. (2019) argued that variability in the along-isobath winds explains 67% of variabil-
 777 ity in the off-shelf transport of fresh water on the southwestern Greenland Shelf. In HY-
 778 COM, the winds, and the resulting Ekman transport, are stronger and show a larger range
 779 of variability when compared to POP. While the POP volume transport varies season-
 780 ally, the HYCOM transport does not display the same seasonality. The difference in trans-
 781 port variability timescale (seasonal in POP and year-to-year in HYCOM) could be one
 782 reason that the 2005-2009 average transports differ between the models. In this region,
 783 surface forcing plays a larger role in influencing cross-shelf transports in POP compared
 784 to HYCOM. In comparing the difference in Ekman transport to the difference in net vol-
 785 ume transport, we find the Southwest region to have the strongest relationship of any
 786 of the sections listed here, with $V_{DW\ diff}$ explaining 5% of the variance in $V_{800\ diff}$.

787 Over the Narrow Shelf (panels B and E), the difference in wind stress between the
 788 two atmospheric reanalysis products is apparent. The transport from surface stress in
 789 POP is small when compared to the cross-isobath volume transport. In both simulations,

790 the seasonal cycle is consistent with strengthening of the along-shelf winds in winter. The
791 stronger onshore transport in HYCOM is consistent with the stronger winds in CFSR
792 (HYCOM, see Figure 6). This difference in transport could be linked to a difference in
793 model forcing or atmosphere-ocean coupling, not model dynamics. The cross-isobath vol-
794 ume transport varies seasonally in POP, but in HYCOM it has a consistent magnitude
795 throughout the year; this is similar to the difference between the two simulations in the
796 Southwest region and consistent with the hypothesis that cross-isobath transport is wind
797 driven.

798 Over the Wide Shelf (panels C and F), the mean and variability of Ekman trans-
799 port are similar between the two forcings. However, the cross-isobath transports have
800 different signs. Note, this is the region in both models where the topographic Rossby waves
801 propagate, with impacts on both volume and heat transport. In POP the variability of
802 the volume transport resembles the Ekman transport, but with opposite sign as has been
803 the case in the other shelf regions. In HYCOM, the transport is consistently onto the
804 shelf, in agreement with the Ekman transport. The topographic Rossby waves travel-
805 ing along the shelf break in this region are likely wind-generated, as discussed in the next
806 section, and influence the volume transport onto the shelf.

807 Overall, to understand where winds drive cross-isobath transport, we look at the
808 regions in both simulations where the transports are correlated to the winds. In the South-
809 west region, along-isobath winds explain similar amounts of variance in POP (6%) and
810 HYCOM (5%). In HYCOM, neither wind forcing explains more than 1% of the variance
811 in transport over the Narrow or Wide Shelf. However, in POP Ekman convergence in
812 the Narrow Shelf region explained 32% of variance in cross-isobath volume transport;
813 no other region had a similar correlation to Ekman convergence. In both simulations,
814 we found that along the Southeast shelf the off-shelf volume transport was the result of
815 relatively weak onshore fluxes being compensated by stronger offshore fluxes at depth.
816 Combined with the low correlation to Ekman transport we conclude that the cross-isobath
817 exchange does not occur primarily in the Ekman layer. These results differ from previ-
818 ous results linking along-shelf winds to off-shelf transport of freshwater in southwest Green-
819 land (Dukhovskoy et al., 2019). We attribute this to the vertical distribution of fresh-
820 water, which is concentrated in the surface layer and directly influenced by Ekman trans-
821 port. In addition, while there are some differences in the wind forcing between simula-
822 tions, the connection between the winds and the transport is not strong enough in ei-



828 **Figure 12.** Plots of spectra computed from daily fields projected net surface stress at every
 829 point along the transect from Denmark Strait to Davis Strait. (A) is results from POP-CORE-
 830 IIAF, (B) is results from HCYOM - CFSR. Vertical dashed lines show the locations of the gates
 831 and solid vertical lines show the region of the propagating signal. In (A) the range of the black
 832 bars is 90 to 500 km for (B) is is from 150 to 270 km, both referenced to the Denmark Strait

823 their simulation to suggest that differences in forcing products are the primary cause of
 824 differences between the two simulations. This is further evidence that Ekman transport
 825 is not a primary driver of cross-shelf transport in these regions.

826 5.2 Mechanism Two: Wind Generation of Topographically Trapped Rossby 827 Waves

833 In the Hovmöller diagrams (Figures 9A and 10A) we observed high-frequency sig-
 834 nals that emanated from a location south of the Denmark Strait. These signals are com-
 835 parable to the topographically trapped Rossby waves (Münchow et al., 2020) in a trough
 836 near the Fram Strait, to the cyclonic eddies formed at the Denmark Strait (Moritz et
 837 al., 2019), and to the coastally trapped shelf waves in this region (Gelderloos et al., 2021).
 838 The period, high phase velocity, and wavelength are more consistent with topographic
 839 Rossby waves than with Kelvin waves or gravity waves. In earlier studies, wind forcing
 840 has been proposed as the primary mechanism that produces similar high frequency sig-
 841 nals (Gelderloos et al., 2021). We focus on the local winds at the shelf break and use fre-
 842 quency spectra (Figure 12) to determine if the frequency of the local winds matches the
 843 frequency of the waves, although the winds do not need to have the frequency of the sig-
 844 nal that they generate. However, this does not exclude the possibility of wind forcing,
 845 Dukhovskoy, Morey, and O'Brien (2009) found tropical storms generated by low-frequency
 846 waves along the Nicaragua Shelf.

847 Spectra are calculated at each point along the shelf break using the winds projected
848 onto the along-isobath direction for each of the three main regions: Southwest, Narrow
849 Shelf, and Wide Shelf. Both wind products show the most variability around 0.05 day^{-1} .
850 The low spatial resolution of both wind products (approximately 90 km for CORE-IAF
851 and approximately 38 km for CFSR) results in the reduced along-shelf variability seen
852 in the spectra; compare to panel C in Figures 9 and 10 where spectra at nearby loca-
853 tions vary. The strong winds over the Narrow Shelf in the HYCOM forcing are likely con-
854 nected to the relatively greater energy in the surface stress near Cape Farewell. At the
855 location along the transect where the signals are generated there is not a clear correspond-
856 ing high-frequency signal in the surface forcing. However, the volume transport at the
857 location of generation is correlated with the surface stress across the region. This pro-
858 vides evidence that the topographic Rossby waves are generated remotely rather than
859 at the shelf break.

860 Comparing the portion of the shelf along which the waves propagate to the spa-
861 tial patterns of EKE in both simulations (Figure 3), we identify regions where the to-
862 pographic Rossby waves may be generated. The elevated EKE extends north of the Den-
863 mark Strait in both simulations to the mouth of Scoresby Sound. In HYCOM, the band
864 of high EKE extends unbroken across the Denmark Strait and is shoreward of the 800 m
865 isobath south of the strait. If we consider the region of elevated EKE to be a proxy for
866 the path of the topographic Rossby waves, this would indicate one reason they are not
867 apparent in HYCOM south of Kangerdlugssuaq Trough is because they are further on-
868 shore than the 800 m bathymetric contour. In POP, the shoreward shifting of the band
869 of elevated EKE occurs further south, consistent with the along-shelf propagation of the
870 topographic Rossby waves. Furthermore, the pattern of elevated EKE indicates these
871 waves could be remotely generated as far north along the shelf as Scoresby Sound. Fur-
872 ther investigation of the generation of these waves is beyond the scope of this work.

873 **6 Conclusion**

874 In order to assess the heat transport onto the Greenland Continental Shelf, we com-
875 pared in detail two high-resolution coupled ocean-sea ice simulations with different at-
876 mospheric forcing. Using a control volume around the 800 m isobath and gates at key
877 locations on the shelf, we determine not only how much heat crosses onto the shelf but
878 also the patterns of transport on the shelf. The region of greatest heat transport onto

879 the shelf is between the Denmark Strait and the Sermilik Troughs in southeast Green-
880 land, where the average heat transport is 16.4 ± 13.8 TW in POP and 55.0 ± 23.3 TW in
881 HYCOM. Currents on the shelf are important in spreading warm water to different shelf
882 regions; in both models the primary source of heat on the southwest continental shelf
883 is from southward transport through the Cape Farewell Gate. The warmest part of the
884 shelf is between the Denmark Strait and the Davis Strait; this region also has the largest
885 seasonal change in heat content.

886 In this study we hypothesised wind-driven transports as one driver of heat trans-
887 port onto the shelf, but volume transport was not clearly linked to off-shore Ekman trans-
888 port. In none of our shelf regions did Ekman dynamics appear to be the primary driver
889 of cross-isobath transport. However, the vertical distribution analyses show that trans-
890 port in the upper layer is compensated by much stronger off-shelf transport at depth.
891 We also quantified how differences in the surface wind stress from our atmospheric forc-
892 ings contributed to the differences in cross shelf transport and found that the wind forc-
893 ings could not explain the difference.

894 Topographic Rossby waves propagate along the shelf break, originating south of
895 the Denmark Strait. These waves have periods of 3-7 days and propagate along the sec-
896 tion of the shelf with the strongest on-shelf volume and heat fluxes. The transport as-
897 sociated with these waves is key in our budget of the Greenland Shelf. They appear to
898 be the leading mechanism of cross shelf exchange and result in a net heat transport onto
899 the shelf. Their occurrence in both simulations indicates that their presence is robust.
900 Similar waves have been observed in other high-resolution simulations. In our simula-
901 tions, local winds do not appear to be a critical driver of wave generation. Further study
902 is needed to identify the key mechanisms including observations that would show evi-
903 dence of these waves.

904 At the outset of this project, our premise for comparing these simulations was that
905 they would provide independent realizations of the Greenland Current system and that
906 the dominant dynamics would be apparent, regardless of model design, if each of the mod-
907 els was reasonably realistic. Both models are similar to observations and to each other
908 in terms of surface currents and EKE, basin temperature and salinity, and volume trans-
909 port through major straits. Fully understanding what difference in the simulation setup
910 results in the difference in dynamics is beyond the scope of this work; however, we hy-

911 pothesize that the primary drivers of differences between the two simulated transports
912 stem from the use of different atmospheric forcings and resolution, and that differing ver-
913 tical coordinate systems could have contributed. Our key findings are consistent between
914 the two simulations despite differences in properties on the shelf and in cross-shelf ex-
915 change.

916 One aspect of the dynamics of the Greenland continental shelf that has been ne-
917 glected in this study is the role of ice sheet meltwater in these cross-shelf exchange mech-
918 anisms. Neither simulation includes a representation of GIS meltwater, which is expected
919 to modify vertical and horizontal shelf stratification. Further simulations are needed to
920 explore the implications of accelerated melting on shelf warming. In addition, our study
921 has shown that mesoscale processes contribute to on-shelf transport. High-resolution stud-
922 ies in this region are needed to understand these processes. Such high-resolution stud-
923 ies could also address the dynamics between the shelf break and the ice sheets that bring
924 the warm water we observe crossing the shelf to the front of glaciers where it drives melt-
925 ing.

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